

High-Rate Quantum Low-Density Parity-Check Codes Assisted by Reliable Qubits

Yuichiro Fujiwara, *Member, IEEE*, Alexander Gruner, *Student Member, IEEE*, and Peter Vandendriessche

Abstract—Quantum error correction is an important building block for reliable quantum information processing. A challenging hurdle in the theory of quantum error correction is that it is significantly more difficult to design error-correcting codes with desirable properties for quantum information processing than for traditional digital communications and computation. A typical obstacle to constructing a variety of strong quantum error-correcting codes is the complicated restrictions imposed on the structure of a code. Recently, promising solutions to this problem have been proposed in quantum information science, where in principle any binary linear code can be turned into a quantum error-correcting code by assuming a small number of reliable quantum bits. This paper studies how best to take advantage of these latest ideas to construct desirable quantum error-correcting codes of very high information rate. Our methods exploit structured high-rate low-density parity-check codes available in the classical domain and provide quantum analogues that inherit their characteristic low decoding complexity and high error correction performance even at moderate code lengths. Our approach to designing high-rate quantum error-correcting codes also allows for making direct use of other major syndrome decoding methods for linear codes, making it possible to deal with a situation where promising quantum analogues of low-density parity-check codes are difficult to find.

Index Terms—Quantum error correction, low-density parity-check code, combinatorial design, entanglement-assisted quantum error-correcting code.

I. INTRODUCTION

QUANTUM error-correcting codes are schemes that recover the original quantum information when the quantum states of quantum bits, or *qubits*, carrying the information are transformed by unintended quantum operations, namely *quantum noise* [1]. As is the case with traditional information processing, it is vital to suppress the effect of quantum noise when processing quantum information. The role of error correction is particularly crucial in the quantum domain because qubits are expected to be highly vulnerable to environmental noise in practical and realistic situations.

While the importance of reliability is apparent, there had been doubts about the existence of a viable scheme for error correction in the quantum domain until the discovery of the

famous 9-qubit code [2] and 7-qubit code [3] in the mid 1990's. These findings ignited intensive and rapidly progressing research on error correction for quantum information. In fact, various types of quantum error-correcting code are now known including the celebrated stabilizer codes [4], [5], which constitute a very general class encompassing the first two quantum error-correcting codes, and codeword stabilized codes [6]. Small quantum error-correcting codes, such as the perfect 5-qubit quantum error-correcting code [7], have been experimentally realized as well [8]–[20].

However, this remarkable progress does not mean that the theory of quantum error correction has become as mature as classical coding theory. It would be more accurate to say that we just started finding ways to realize quantum error correction while cleverly circumventing challenging obstacles imposed by quantum mechanical phenomena.

For instance, while the stabilizer formalism developed in [21] has given rise to a wide range of quantum error-correcting codes, one of the theoretically challenging problems with this approach is that the admissible structures of a code are severely restricted when compared to the freedom we have in classical code design. The fact that there are only few successful general frameworks for quantum code design also limits the variety of quantum error-correcting codes, which is a crucial problem because actual realizations of large-scale quantum information processing is expected to demand various types of peculiar requirement.

One effective way to overcome the limitations and difficulties in the quantum domain is to develop a fresh and quantum mechanically valid framework that makes it possible to directly import a wider range of classical coding theory to the quantum regime. The *entanglement-assisted stabilizer formalism* is a major breakthrough in this direction, where one may fully exploit any binary or quaternary linear code over the binary field \mathbb{F}_2 or the finite field \mathbb{F}_4 of order four respectively for correcting errors on qubits as long as there is an adequate supply of maximally entangled noiseless qubits, called *ebits*, to assist quantum error correction [22]. Entanglement-assisted quantum error-correcting codes can be regarded as generalized stabilizer codes in that those requiring no ebit are exactly the standard stabilizer codes; if a linear code can not be turned into a quantum error-correcting code through the standard stabilizer formalism, one may still exploit it by assuming that some amount of quantum information can be shared through a noiseless channel as ebits to help encode and decode noisy qubits.

A major drawback of entanglement assistance is that completely noiseless qubits are extremely difficult to provide in

This work was supported by JSPS (Y.F.), LGFG Baden-Württemberg (A.G.), and FWO (P.V.). The third author is support by a PhD fellowship of the Research Foundation - Flanders (FWO).

Y. Fujiwara is with the Division of Physics, Mathematics and Astronomy, California Institute of Technology, MC 253-37, Pasadena, CA 91125 USA (email: yuichiro.fujiwara@caltech.edu).

A. Gruner is with the Wilhelm Schickard Institute for Computer Science, Eberhard Karls Universität Tübingen, Sand 13, D-72076 Tübingen, Germany (email: alexander.gruner@uni-tuebingen.de).

P. Vandendriessche is with the Department of Mathematics, Ghent University, Krijgslaan 281 - S22, 9000 Ghent, Belgium (email: pv@cage.ugent.be).

a practical quantum device. This disadvantage is particularly pronounced in the context of storing quantum information, where the information source and sink may not be spatially distant but are separate in the time domain. This characteristic of entanglement-assisted quantum error-correcting codes led to a series of research trying to identify excellent linear codes which can be imported by relying only on a tiny number of ebits [23]–[34]. Playing a crucial role in these theoretical results is the assumed future technology of manipulating a small number of qubits with extreme reliability to realize perfect and stable ebits.

Very recently, a framework that significantly reduces the burden of providing extreme reliability has been proposed, where any binary or quaternary linear codes over \mathbb{F}_2 or \mathbb{F}_4 respectively can be fully exploited as long as we can provide auxiliary qubits that are only subject to a restricted quantum error model [35]. This framework takes advantage of the fact that while realizing completely noiseless qubits is a very difficult task, not every kind of quantum error is equally difficult to suppress through technical development on hardware. For instance, it is known that one can correct any type of quantum error in the standard general error model if two particular types of error, called a bit error and phase error, can be corrected under the assumption that both may happen on the same qubit [1]. However, phase errors due to dephasing are expected to be far more likely than bit errors in most actual quantum devices [36], which implies that bit errors would be far less problematic. In the newer framework, one may choose an error model in such a way that most qubits can suffer from bit errors and phase errors while only phase errors may occur on the auxiliary qubits. Hence, unlike the entanglement-assisted stabilizer formalism, which requires completely noiseless auxiliary qubits, the newer framework only needs more easily achievable “less noisy” ones.

With all these advances in this field, one may think that the problem of severely restricted structures of quantum error-correcting codes is largely solved. The caveat is that the statement that a given linear code \mathcal{C} can be imported as a stabilizer code, entanglement-assisted quantum error-correcting code or one that is assisted by less noisy qubits only means that \mathcal{C} admits a suitable parity-check matrix that is exploitable for quantum error correction. In other words, of all distinct parity-check matrices that define one same linear code \mathcal{C} , only some special ones are usable.

To illustrate this problem, consider linear codes that greatly benefit from parity-check matrices in particular form in the classical domain. For instance, *low-density parity-check* (LDPC) codes are linear codes that admit parity-check matrices with a small number of nonzero entries such that iterative decoding performs well [37]. They are among the state-of-the-art error-correcting codes in classical coding theory in the sense that well-designed LDPC codes almost achieve the Shannon limit and have remarkably low decoding complexity. Since we have a means to import the theory of linear codes into the quantum domain, LDPC codes constitute very promising ingredients for quantum error correction. However, whether a given linear code is qualified as an excellent LDPC code depends on whether it has a parity-check matrix suitable for

iterative decoding. This implies that whether we may have a quantum counterpart that inherits the attractive characteristics of a given LDPC code depends on whether its particular parity-check matrix suitable for iterative decoding is compatible with the chosen method for turning a linear code into a quantum error-correcting code.

The purpose of the present paper is to give insight into how best to exploit the recently proposed frameworks for quantum error correction assisted by reliable qubits when the restrictions on parity-check matrices must be taken into account. In particular, we mostly focus on the case when excellent LDPC codes are used to simultaneously achieve very low decoding complexity and high error correction performance. To take full advantage of auxiliary qubits while keeping our work well-focused, we aim to construct quantum error-correcting codes with a few other properties that would be desirable in various situations.

An $[[n, k]]$ quantum error-correcting code of *length* n and *dimension* k encodes k -qubit information into n physical qubits, where the two nonnegative integers n and k satisfy the condition that $n > k \geq 0$. The first property we aim for is a very high information rate in the absolute sense, which means that we would like an $[[n, k]]$ quantum error-correcting code with k close to n . In addition to this condition, we strive to restrict ourselves to quantum error-correcting codes of modest and realistic length. Thus, we do not consider the case when the parameter n is an unrealistically large integer or the purely theoretical case of n approaching infinity. To illuminate the potential of our approach, we aim for simultaneously satisfying these demanding conditions while achieving high error correction performance comparable to what would be attainable in a hypothetical situation where some of the best known classical LDPC codes were freely available for quantum error correction without the limitation on the structure of parity-check matrices. As we will see later, carefully designed quantum LDPC codes can achieve this goal through assisted quantum error correction. Furthermore, a brief discussion at the end of this paper will show how assisted quantum error correction with less noisy qubits, if exploited with a different decoding method for linear codes, may remain successful in a situation where excellent quantum LDPC codes are difficult to construct.

In the next section we briefly review quantum error-correcting codes assisted by reliable qubits. Section III discusses the use of LDPC codes of high rate for quantum error correction through the recently proposed frameworks. We examine the performance of our quantum LDPC codes through simulations in Section IV. Concluding remarks including a brief discussion on how to apply assisted quantum error correction to other decoding methods are given in Section V.

II. QUANTUM ERROR CORRECTION WITH RELIABLE AUXILIARY QUBITS

In this section we give a brief review of how reliable auxiliary qubits help correct quantum errors. For the basics of quantum information theory, we refer the reader to [1]. All facts in classical coding theory we use in this section can be found in [38].

As usual, by a binary linear $[n, k, d]$ code, we mean a k -dimensional subspace \mathcal{C} of the n -dimensional vector space over \mathbb{F}_2 in which a nonzero vector with the smallest number of nonzero entries has exactly d nonzero entries, that is, $\min\{\text{wt}(c) \mid c \in \mathcal{C}, c \neq 0\} = d$. Because we only consider a binary code, we omit the term binary when referring to linear codes and LDPC codes. As stated earlier, an $[[n, k]]$ quantum error-correcting code encodes k logical qubits into n physical qubits, which is analogous to a linear $[n, k, d]$ code in the sense that the classical code encodes k logical bits into n physical bits.

An important fact in the quantum domain is that, through a process called *discretization*, an error correction scheme can correct any general quantum error on one qubit if it can correct the effects of the Pauli operators X, Z and their product XZ , where the operator X corresponds to the *bit error* on one qubit while Z represents the *phase error* [1]. Similarly, quantum errors on multiple qubits can be corrected if the corresponding transformation by a combination of X, Z and both at the same time on each of the affected qubits is detected and reversed.

The quantum error-correcting codes we consider in this paper also take advantage of discretization. Hence, without loss of generality, we always assume that a quantum channel may introduce on each qubit only a bit error, a phase error or both at the same time as a quantum error during information transmission unless otherwise stated.

The rest of this section is divided into two subsections. Section II-A presents the basics of the framework for quantum error correction given in [35] that is assisted by qubits on which only one particular kind of quantum error may occur. We briefly review in Section II-B the entanglement-assisted stabilizer formalism developed in [22] which uses completely noiseless qubits.

A. Less Noisy Auxiliary Qubits

Here we give the basics of quantum error correction assisted by less noisy qubits from the viewpoint of classical coding theory. The following is the tool we use to import linear codes.

Theorem 2.1 ([35]): If there exists a linear $[n, k, d]$ code, then there exist unitary operations that encode k logical qubits into $2n - k$ physical qubits and correct up to $\lfloor \frac{d-1}{2} \rfloor$ quantum errors under the assumption that a fixed set of $2(n-k)$ physical qubits may experience phase errors but no bit errors.

Roughly speaking, the above theorem says that any linear $[n, k, d]$ code, which corrects errors on up to $\lfloor \frac{d-1}{2} \rfloor$ bits, can be turned into a $[[2n - k, k]]$ quantum error-correcting code that corrects quantum errors on up to $\lfloor \frac{d-1}{2} \rfloor$ qubits as long as predetermined $2(n-k)$ qubits are only subject to phase errors. Note that if the original linear $[n, k, d]$ code is of sufficiently high rate, the $2(n-k)$ auxiliary qubits consist of only a small fraction of the $2n - k$ physical qubits.

A particularly useful fact regarding this type of quantum error correction is that we can employ decoding methods for linear codes based on error syndromes. We formulate this most fundamental part of our approach in the form of a theorem below.

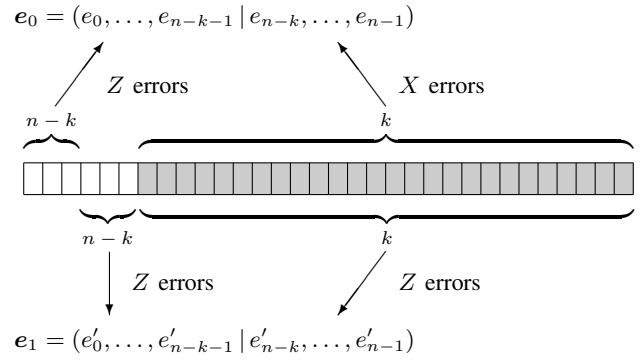


Fig. 1. Correspondence of quantum errors to error vectors. | The white boxes represent the $2(n-k)$ less noisy qubits that may experience only phase errors. The gray boxes are the k noisy qubits that may suffer from bit and/or phase errors. The first $n-k$ bits of e_0 and the first $n-k$ bits of e_1 correspond to whether phase errors occurred on the $2(n-k)$ less noisy qubits. The remaining k bits of e_0 indicate whether bit errors occurred on the k noisy qubits while the remaining k bits of e_1 correspond to phase errors on these noisy qubits.

Theorem 2.2: Let \mathcal{C} be a linear $[n, k, d]$ code. Assume that $2n - k$ physical qubits q_i , $0 \leq i \leq 2n - k - 1$, are sent through a noisy quantum channel in which the first $2(n-k)$ qubits q_i , $0 \leq i \leq 2(n-k) - 1$ are only subject to phase errors while the remaining k qubits q_i , $2(n-k) \leq i \leq 2n - k - 1$ are subject to both bit errors and phase errors. Define a pair $e_0 = (e_0, \dots, e_{n-1}), e_1 = (e'_0, \dots, e'_{n-1}) \in \mathbb{F}_2^n$ of n -dimensional vectors such that for $0 \leq i \leq n-k-1$, $e_i = 1$ if a phase error occurred on q_i and $e_i = 0$ otherwise, such that for $n-k \leq i \leq n-1$, $e_i = 1$ if a bit error occurred on q_{i+n-k} and $e_i = 0$ otherwise, and such that for $0 \leq i \leq n-1$, $e'_i = 1$ if a phase error occurred on q_{i+n-k} and $e'_i = 0$ otherwise. Let H be a parity-check matrix of \mathcal{C} in standard form. There exists a $[[2n - k, k]]$ quantum error-correcting code that allows for retrieving classical information about quantum errors in the form of a pair $s_0, s_1 \in \mathbb{F}_2^{n-k}$ of $(n-k)$ -dimensional vectors such that $s_0 = He_0^T$ and $s_1 = He_1^T$.

Note that the binary vectors e_0 and e_1 in the above theorem specify what type of quantum error occurred on which qubit. The correspondence between each bit of the error vectors e_0, e_1 and the type and location of each quantum error is summarized in Fig. 1. The point of Theorem 2.2 is that because H is a parity-check matrix of a linear code of minimum distance d , we can correctly infer e_0 and e_1 from the syndromes s_0 and s_1 , which are He_0^T and He_1^T respectively, if the weights of e_0 and e_1 are both less than or equal to $\lfloor \frac{d-1}{2} \rfloor$. This implies that, by the definition of e_0 and e_1 , the positions of all bit errors and phase errors can be identified and thus corrected if the number of physical qubits that suffer bit errors, phase errors or both is at most $\lfloor \frac{d-1}{2} \rfloor$.

While it is straightforward to derive Theorem 2.2 from the results already presented in [35], for completeness, we give a formal proof in the remainder of this subsection.

For a unitary operator U and a v -dimensional vector $\mathbf{a} = (a_0, \dots, a_{v-1}) \in \mathbb{F}_2^v$, define $U^{\mathbf{a}}$ as the v -fold tensor product $O_0 \otimes \dots \otimes O_{v-1}$, where $O_i = U$ if $a_i = 1$ and O_i is the identity operator otherwise.

Take a linear $[n, k, d]$ code with a parity-check matrix H in standard form

$$H = \begin{bmatrix} I & A \end{bmatrix}$$

for some $(n-k) \times k$ matrix A over \mathbb{F}_2 , where I is the $(n-k) \times (n-k)$ identity matrix. The Z -information check matrix H_Z and X -information check matrix H_X of H are the $2(n-k) \times k$ matrices

$$H_Z = \begin{bmatrix} A \\ 0 \end{bmatrix}$$

and

$$H_X = \begin{bmatrix} 0 \\ A \end{bmatrix}$$

respectively. Simply put, H_Z and H_X are matrices composed of the $(n-k) \times k$ all-zero matrix and the columns of the parity-check matrix H that correspond to the information bits.

Let $|0\rangle_X^{\otimes 2(n-k)}$ be $2(n-k)$ qubits in the joint $+1$ eigenstate of $X^{\otimes 2(n-k)}$. Without loss of generality, we assume that $|0\rangle_X = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$ and that $|1\rangle_X = \frac{|0\rangle - |1\rangle}{\sqrt{2}}$, where $|0\rangle$ and $|1\rangle$ are the computational basis.

Lemma 2.3 ([35]): Assume that there exists a linear code of length n and dimension k with a parity-check matrix H in standard form. Define

$$Q = \sum_{\mu \in \mathbb{F}_2^{2(n-k)}} |\mu\rangle \langle \mu| \otimes X^{\mu H_X} Z^{\mu H_Z}$$

and Q^\dagger as its complex conjugate, where H_Z and H_X are the Z - and X -information check matrices of H . Take a pair $e_X, e_Z \in \mathbb{F}_2^{2n-k}$ of arbitrary $(2n-k)$ -dimensional vectors. Define e_{X_l} and e_{X_r} as the first $2(n-k)$ and the remaining k bits of e_X respectively so that $e_X = (e_{X_l}, e_{X_r})$. Define similarly $e_Z = (e_{Z_l}, e_{Z_l}, e_{Z_r})$, where e_{Z_l} , e_{Z_l} , and e_{Z_r} are the first $n-k$, the next $n-k$, and the last k bits of e_Z respectively. Let $e_0 = (e_{Z_l}, e_{X_r})$ and $e_1 = (e_{Z_l}, e_{Z_r})$. Assume that $e_{X_l} = 0$. For arbitrary k qubit state $|\psi\rangle$,

$$\begin{aligned} Q^\dagger X^{e_X} Z^{e_Z} Q |0\rangle_X^{\otimes 2(n-k)} |\psi\rangle \\ = |(He_0^T, He_1^T)\rangle_X \otimes X^{e_{X_r}} Z^{e_{Z_r}} |\psi\rangle. \end{aligned} \quad (1)$$

Theorem 2.2 immediately follows from the above lemma.

Proof of Theorem 2.2: Regard the arbitrary k qubit state $|\psi\rangle$, unitary operator Q , and complex conjugate Q^\dagger in Lemma 2.3 as the original k -qubit information which is to be encoded, an encoding operator, and a decoding operator respectively. Assume that the supports $\text{supp}(e_X), \text{supp}(e_Z)$ of the pair e_X, e_Z of arbitrary $(2n-k)$ -dimensional vectors represent the positions of X errors and Z errors introduced by a quantum channel respectively such that an X error occurred on the i th physical qubit if and only if $i \in \text{supp}(e_X)$ and such that a Z error occurred on the i th physical qubit if and only if $i \in \text{supp}(e_Z)$. With these assumptions, the other assumption that $e_{X_l} = 0$ made in Lemma 2.3 corresponds to the condition that a fixed $2(n-k)$ physical qubits are only subject to phase errors. Measuring the $2(n-k)$ ancilla qubits on the right-hand side of Equation (1) in the Hadamard rotated basis gives a $2k$ -dimensional vector $s \in \mathbb{F}_2^{2k}$ of which the first half is the k -dimensional vector $s_0 = He_0^T$ and the second half of which is $s_1 = He_1^T$. The proof is complete. ■

Clearly, if $|\text{supp}(e_0)|, |\text{supp}(e_1)| \leq \lfloor \frac{d-1}{2} \rfloor$, the retrieved classical information in the form of a pair of k -dimensional vectors He_0^T, He_1^T uniquely identifies the locations of nonzero bits in e_0 and e_1 as in standard syndrome decoding for linear codes. The assumption that the number of quantum errors is $\lfloor \frac{d-1}{2} \rfloor$ or less, which means $|\text{supp}(e_X) \cup \text{supp}(e_Z)| \leq \lfloor \frac{d-1}{2} \rfloor$, implies that both $|\text{supp}(e_0)|$ and $|\text{supp}(e_1)|$ are less than or equal to $\lfloor \frac{d-1}{2} \rfloor$. Trivially, once e_0 and e_1 are correctly inferred, the two vectors e_X and e_Z that specify the positions of bit errors and phase errors can be fully reconstructed.

B. Entanglement Assistance

Here we review necessary basic facts on entanglement-assisted quantum error-correcting codes. We follow the method used in [28], [31] for constructing quantum LDPC codes through the entanglement-assisted analogue of the Calderbank-Shor-Steane (CSS) construction [3], [39]. Similar to the previous method that uses less noisy qubits, this entanglement-assisted method allows for extracting the information about what type of quantum error occurred on which qubit by simply treating a pair of binary vectors as error syndromes.

An $[[n, k; c]]$ *entanglement-assisted quantum error-correcting code* is an error correction scheme that encodes k logical qubits into n physical qubits with the help of c ebits. The c ebits are sent through a noiseless channel. When importing a linear code with a parity-check matrix H , the required number c of noiseless qubits is exactly the 2-rank of the matrix HH^T over \mathbb{F}_2 [24]. Because we only consider ranks over \mathbb{F}_2 , in what follows we simply write $\text{rank}(A)$ to mean the 2-rank of a given matrix A .

The following is a straightforward consequence of the CSS construction for entanglement-assisted quantum error-correcting codes.

Theorem 2.4 ([31]): Assume that n physical qubits are sent through a noisy quantum channel. Define $e_X = (e_0, \dots, e_{n-1}) \in \mathbb{F}_2^n$ to be the n -dimensional vector representing the positions of bit errors such that $e_i = 1$ if a bit error occurred on the i th qubit and $e_i = 0$ otherwise. Define also $e_Z = (e'_0, \dots, e'_{n-1}) \in \mathbb{F}_2^n$ to be the n -dimensional vector representing the positions of phase errors such that $e'_i = 1$ if a phase error occurred on the i th qubit and $e'_i = 0$ otherwise. Let H be a parity-check matrix of a linear $[n, k, d]$ code. There exists an $[[n, 2k - n + \text{rank}(HH^T); \text{rank}(HH^T)]]$ entanglement-assisted quantum error-correcting code that allows for retrieving classical information about quantum errors in the form of a pair $s_X, s_Z \in \mathbb{F}_2^{n-k}$ of $(n-k)$ -dimensional vectors such that $s_X = He_X^T$ and $s_Z = He_Z^T$.

Because H is a parity-check matrix of a linear $[n, k, d]$ code, it is straightforward to see that if the number of bit errors and that of phase errors are both less than or equal to $\lfloor \frac{d-1}{2} \rfloor$, the qubits on which bit errors occurred and the ones on which phase errors occurred can be identified. Note that unlike in Theorem 2.2, we do not require H to be in standard form or full-rank in Theorem 2.4. Instead, typical and realistic assumptions require that the 2-rank $\text{rank}(HH^T)$ be very small because it is the number of ebits we need to engineer extremely accurately and protect perfectly. The most extreme case is

when $\text{rank}(HH^T) = 0$, where Theorem 2.4 reduces to the CSS construction in its original form. Entanglement assistance takes place when $\text{rank}(HH^T) \geq 1$.

III. ASSISTED QUANTUM LDPC CODES

In this section we study the desirable structures of parity-check matrices for use in high-rate quantum error correction assisted by less noisy qubits or error-free ebits. We make the conservative assumption that the receiver has no knowledge of possible correlations between bit errors and phase errors so that the decoder approximates the quantum channel by two binary symmetric channels, one of which introduces the operator X independently on each physical qubit with probability p_x and the other of which make the operator Z act independently on each physical qubit with probability p_z . Thus, in the case of codes assisted by less noisy auxiliary qubits, the receiver employs two separate decoders for a linear code to infer e_0 and e_1 in Theorem 2.2 under the condition that the parity-check matrix H and two binary vectors $s_0 = He_0^T$, $s_1 = He_1^T$ are given. For entanglement-assisted quantum error-correcting codes, two separated decoders are used to infer e_X and e_Z in Theorem 2.4 from two binary vectors $s_X = He_X^T$, $s_Z = He_Z^T$ and the parity-check matrix H . In both cases, the receiver employs the sum-product algorithm for inference [37].

It is notable that if the receiver has some knowledge of correlations between bit errors and phase errors, this information can be incorporated into the decoding algorithm with an increase in decoding complexity by carefully implementing a quantum analogue of brief propagation [40]. In fact, significant improvements in error correction performance have been reported in a very optimistic situation where the receiver has perfect knowledge of a channel with a very strong correlation due to depolarizing noise [41]–[43]. Compared to this ideal assumption, our setting assumes a smaller amount of exploitable information about the channel. This allows us to give a conservative estimate on error correction performance as a likely lower bound for various situations and avoid the risk of relying on unrealistically accurate knowledge of how quantum errors manifest on actual hardware.

We divide the remainder of this section into three subsections. Section III-A provides the definitions of combinatorial designs we take advantage of for designing codes throughout this paper. In Section III-B we study parity-check matrices suitable for use as quantum LDPC codes assisted by less noisy auxiliary qubits. Desirable parity-check matrices for entanglement-assisted quantum LDPC codes are investigated in Section III-C.

A. Combinatorial Designs

Let K be a subset of positive integers. A *pairwise balanced design of order v and index 1 with block sizes from K* , denoted by $\text{PBD}(v, K, 1)$, is an ordered pair (V, \mathcal{B}) , where V is a nonempty finite set of v elements, called *points*, and \mathcal{B} is a set of subsets of V , called *blocks*, that satisfies the following two conditions:

- (i) each unordered pair of distinct elements of V appears in exactly one block of \mathcal{B} ,
- (ii) for every $B \in \mathcal{B}$ the cardinality $|B| \in K$.

When K is a singleton $\{\mu\}$, the PBD is a *Steiner 2-design of order v and block size μ* , and is denoted by $S(2, \mu, v)$. A simple counting argument shows that the number of blocks in an $S(2, \mu, v)$ is exactly $\frac{v(v-1)}{\mu(\mu-1)}$. A PBD of order v is *trivial* if it has no blocks or consists of only one block of size v . The trivial PBD with no blocks necessarily has only one point.

Define $\alpha(K) = \gcd\{\mu - 1 \mid \mu \in K\}$ and $\beta(K) = \gcd\{\mu(\mu - 1) \mid \mu \in K\}$. Necessary conditions for the existence of a $\text{PBD}(v, K, 1)$ are $v - 1 \equiv 0 \pmod{\alpha(K)}$ and $v(v - 1) \equiv 0 \pmod{\beta(K)}$ [44]. These conditions are known to be asymptotically sufficient.

Theorem 3.1 (Wilson [45]): There exists a constant v_K such that for every $v > v_K$ satisfying $v - 1 \equiv 0 \pmod{\alpha(K)}$ and $v(v - 1) \equiv 0 \pmod{\beta(K)}$ there exists a $\text{PBD}(v, K, 1)$.

An *incidence matrix* of a PBD (V, \mathcal{B}) with $|V| = v$ and $|\mathcal{B}| = b$ is a binary $v \times b$ matrix $H = (h_{i,j})$ with rows indexed by points, columns indexed by blocks, and $h_{i,j} = 1$ if the i th point is contained in the j th block, and $h_{i,j} = 0$ otherwise.

It is known that incidence matrices of PBDs of index 1 are generally good candidates of parity-check matrices of LDPC codes for high speed information transmission because of their good error tolerance at relatively short lengths [46]–[49]. Our goal in the following two subsections is to identify and give explicit constructions for particularly promising classes of PBDs whose incidence matrices may be used as parity-check matrices for assisted quantum error correction.

B. Parity-Check Matrices for Phase Error Qubit Assistance

The explicit restriction on the structure of a parity-check matrix H of a linear $[n, k, d]$ code in Theorem 2.2 is that it must be in standard form

$$H = \begin{bmatrix} I & A \end{bmatrix}$$

for some $(n-k) \times k$ matrix A over \mathbb{F}_2 , where I is the $(n-k) \times (n-k)$ identity matrix. As we will see later in this subsection, incidence matrices of PBDs of index 1 may be seen as parity-check matrices that give the largest possible information rates among all possible H avoiding certain undesirable structures for the standard sum-product algorithm. Because our goal is to construct promising parity-check matrices of LDPC codes of extremely high rate, here we would like H as a whole to form an incidence matrix of a PBD of index 1. The following proposition allows us to only consider the part A in this regard.

Proposition 3.2: Let $H = \begin{bmatrix} I & A \end{bmatrix}$ be a parity-check matrix of a linear code of length n and dimension k in standard form. H is an incidence matrix of a PBD of index 1 if and only if the $(n-k) \times k$ matrix A is an incidence matrix of a PBD of index 1 containing no block of size 1.

Proof: Assume that H is a parity-check matrix in standard form that forms an incidence matrix of a PBD of index 1. By definition, the PBD contains exactly $n - k$ blocks of size 1, which correspond to the $n - k$ columns of weight 1 in H . Because these blocks do not contribute to the number of each pair of points appearing in blocks, deleting the corresponding

$(n-k) \times (n-k)$ identity matrix I leaves $(n-k) \times k$ matrix A , where every pair of rows have exactly one column in which both rows have 1. Indexing rows by points and columns by blocks, A forms an incidence matrix of a PBD of index 1. Because every column of weight 1 is deleted from H , this PBD does not have a singleton as a block. Conversely, because a block of size 1 does not have a pair of points, combining the $(n-k) \times (n-k)$ identity matrix I and an incidence matrix A of a PBD($n-k, K, 1$) with $1 \notin K$ gives an incidence matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ of a PBD of index 1. ■

In view of the above proposition, we would like to find incidence matrices A of PBDs without blocks of size 1 that do not contain or produce harmful structures when combined with the identity matrices to obtain valid parity-check matrices $H = \begin{bmatrix} I & A \end{bmatrix}$ for Theorem 2.2. While it is generally a very difficult open problem to exactly determine relative harmfulness of each substructure of a parity-check matrix for the sum-product algorithm over a binary symmetric channel, there are known structures that have theoretically or empirically been shown to be undesirable (see [50] and references therein). We first consider a few of the more harmful structures.

The *Tanner graph* of an $m \times n$ parity-check matrix H is the bipartite graph consisting of n *bit vertices* indexed by bits of the corresponding code and m *parity-check vertices* indexed by parity-check equations defined by H , where an edge joins a bit vertex to a parity-check vertex if the bit is involved in the corresponding parity-check equation. An l -cycle in a graph is a sequence of $l + 1$ connected vertices which starts and ends at the same vertex in the graph and contains no other vertices more than once. Clearly, a 4-cycle in a Tanner graph is equivalent to a 2×2 all-one submatrix in a parity-check matrix. A 6-cycle is a 3×3 submatrix in which each row and column has exactly two ones. The *girth* of a parity-check matrix is the length of a shortest cycle in the corresponding Tanner graph. Since a Tanner graph is bipartite, its girth is always even. When it is clear from context which parity-check matrix is considered, we may speak of the “girth of an LDPC code.” It is known that very short cycles tend to be harmful when the sum-product algorithm is employed. In particular, 4-cycles have a very noticeable negative effect on the error correction performance of the sum-product algorithm [51]. For this reason, we would like the girth of a parity-check matrix to be strictly larger than 4.

Because an LDPC code is a linear code equipped with a particular decoding algorithm, the minimum distance also plays a role. While the sum-product algorithm is generally less sensitive to the minimum distance than other simple decoding methods, this fundamental parameter is especially important to a code of very high rate because its very large dimension dictates that the minimum distance be small compared to the length. The following proposition concerns with the number of short cycles and minimum distances of parity-check matrices based on incidence matrices of PBDs together with columns of weight 1.

Proposition 3.3: Let A be an $(n-k) \times k$ incidence matrix of a nontrivial PBD($n-k, K, 1$) with $1 \notin K$. Then the binary matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ is a parity-check matrix of a linear $[n, k, d]$ code in standard form whose girth is 6 and minimum

distance $d = 1 + \min\{\mu \mid \mu \in K\}$.

Proof: Because no pair of points appear twice in a PBD of index 1, there exists no 2×2 all-one submatrix in A . Hence, the girth of A is at least 6. Take an arbitrary column c_1 of A . Write the block B_1 which corresponds to c_1 as $\{v_1, \dots, v_{|B_1|}\}$. Because the PBD is nontrivial, every row of A has at least two ones. Thus, there exists a column c_2 which corresponds to another block $B_2 = \{v_1, v_{|B_1|+1}, \dots, v_{|B_1|+|B_2|-1}\}$, where $v_i \neq v_j$ for any i and j , $i \neq j$. Take the column c_3 representing the block B_3 that contains the pair $\{v_2, v_{|B_1|+1}\}$. The three columns c_1 , c_2 , and c_3 induce a 6-cycle, which implies that the girth of A is exactly 6. Since joining the identity matrix I does not introduce 4-cycles, the girth of H is exactly 6. It suffices to show that a smallest set of linearly dependent columns in H is of cardinality $1 + \min\{\mu \mid \mu \in K\}$. Because no pair of points appear twice in a PBD of index 1, any set of linearly dependent columns that contains at least one column of A is of cardinality at least $1 + \min\{\mu \mid \mu \in K\}$. All columns in the identity matrix are linearly independent. Take an arbitrary column c of A whose weight is the smallest. The identity matrix I contains the set S of columns of cardinality $\min\{\mu \mid \mu \in K\}$ such that $S \cup \{c\}$ forms a set of linearly dependent columns. Because no pairs appear twice in a PBD of index 1, any set of columns containing more than one column from A is linearly independent if its cardinality is less than or equal to $\min\{\mu \mid \mu \in K\}$. ■

As we stated earlier in this subsection, incidence matrices of PBDs of index 1 are extreme in terms of information rates in the following sense.

Proposition 3.4: Let H be a parity-check matrix in standard form that forms an incidence matrix of a nontrivial PBD of order $n-k$ and index 1. Any parity-check matrix of the same size, same column weight distribution, same or higher girth, and same or higher minimum distance as H is of the same or lower rate as H .

Proof: Let H be an $(n-k) \times n$ parity-check matrix in standard form and $w_c = (w_0, w_1, \dots, w_{n-1})$ an n -dimensional vector over nonnegative integers \mathbb{N}_0 such that the column $c_i = (c_0, c_1, \dots, c_{n-k})$ of $H = (c_0, \dots, c_{n-1})$ contains exactly w_i 1s. Without loss of generality, we assume that $w_i = 1$ for $0 \leq i \leq n-k-1$. If there exists i such that $n-k \leq i \leq n-1$ and $w_i = 0$ or 1, then the minimum distance of the corresponding linear code is less than or equal to 2. Assume that $w_i \neq 0, 1$ for $n-k \leq i \leq n-1$. If H is an incidence matrix of a PBD of index 1, then by counting the number of pairs of points in blocks, we have

$$\sum_{i=n-k}^{n-1} \binom{w_i}{2} = \binom{n-k}{2}.$$

Note that a 2×2 all-one submatrix corresponds to the same pair appearing more than once and induces a 4-cycle. Thus, the above equation dictates that any parity-check matrix of the same size and same column weight distribution w_c is either of girth 4 or an incidence matrix of a PBD of index 1. ■

Proposition 3.5: Let H be a parity-check matrix in standard form that forms an incidence matrix of a nontrivial PBD of order $n-k$ and index 1. Take a nonzero vector $c \in \mathbb{F}_2^{n-k}$ of

dimension $n - k$. Adding c^T to H as a column results in a parity-check matrix of an LDPC code of minimum distance 2 or girth 4.

Proof: If c is of weight 1, then one of the columns in the $(n - k) \times (n - k)$ identity matrix is identical to c , resulting in a parity-check matrix of a linear code of minimum distance 2. If c is of weight 2 or larger, then the added column c^T has a pair r, r' of rows whose entries are both 1. Because every pair of points appear exactly once in a PBD of index 1, its incidence matrix H contains a column where the entries of the pair r, r' of rows are both 1. Thus, adding c^T to H introduces a 2×2 submatrix. ■

Because a parity-check matrix H in standard form is always of full rank, roughly speaking, the above propositions state that H forming an incidence matrix of a PBD of index 1 achieves the largest possible number of information bits among all parity-check matrices in standard form with the same number of check equations and the same degree distribution under the constraint that the associated LDPC code must be of minimum distance larger than 2 and girth greater than 4.

An incidence matrix A of a $\text{PBD}(n - k, K, 1)$ gives an LDPC code of minimum distance $1 + \min\{\mu \mid \mu \in K\}$ when combined with the identity matrix. Hence, increasing the smallest block size improves the minimum distance. However, because a block of size x contains $\binom{x}{2}$ pairs, a block of larger size contains more pairs of points. Since avoiding 4-cycles while achieving the highest possible rate is equivalent to packing as many different pairs of points as possible in a set of blocks while including no pair of points more than once, increasing block sizes lowers the achievable information rate in general. Hence, we consider the case when the column weights of the matrix A are uniform. This means that K is a singleton $\{\mu\}$, that is, the corresponding $\text{PBD}(n - k, K, 1)$ is a Steiner 2-design $S(2, \mu, n - k)$. As stated earlier in Section III-A, an $S(2, \mu, v)$ contains exactly $\frac{v(v-1)}{\mu(\mu-1)}$ blocks. Thus, the corresponding code dimension can be quite large at a moderate length.

Proposition 3.6: Let A be an incidence matrix of an $S(2, \mu, v)$ and I a $v \times v$ identity matrix. A parity-check matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ defines an LDPC code of length $\frac{v(v-1)}{\mu(\mu-1)} + v$, dimension $\frac{v(v-1)}{\mu(\mu-1)}$, girth 6, and minimum distance $\mu + 1$.

As can be seen from the above proposition, the rates of LDPC codes defined by incidence matrices of $S(2, \mu, v)$ s become close to 1 very quickly as v tends to infinity. Theorems 2.1 and 2.2 assure that the corresponding quantum error-correcting codes assisted by less noisy qubits inherit this characteristic.

Theorem 3.7: Let A be an incidence matrix of an $S(2, \mu, v)$ and I a $v \times v$ identity matrix. There exists a $\left\lceil \frac{v(v-1)}{\mu(\mu-1)} + 2v, \frac{v(v-1)}{\mu(\mu-1)} \right\rceil$ quantum error-correcting code that identifies the types and locations of quantum errors through the LDPC code defined by the parity-check matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ under the assumption that a fixed set of $2(n - k)$ physical qubits may experience phase errors but no bit errors.

One strategy to improve the error correction performance under the sum-product algorithm is to decrease the number of structures that are responsible for dominating errors. While

joining the identity matrix and the incidence matrix A of a Steiner 2-design of block size μ always results in an LDPC code of minimum distance $\mu + 1$, it is desirable for the linear code defined by A alone to have a larger minimum distance because it eliminates dominating sources of errors to an extent.

It is trivial that the minimum distance of a linear code whose parity-check matrix forms an incidence matrix of an $S(2, \mu, v)$ is at least $\mu + 1$. To investigate the minimum distances of linear codes based on Steiner 2-designs further, we define some combinatorial design theoretic notions. A *configuration* \mathcal{C} in an $S(2, \mu, k)$, (V, \mathcal{B}) , is a subset $\mathcal{C} \subseteq \mathcal{B}$ of the block set. The set of points appearing in at least one block of a configuration \mathcal{C} is denoted by $V(\mathcal{C})$. Two configurations \mathcal{C} and \mathcal{C}' are *isomorphic* if there exists a bijection $\phi : V(\mathcal{C}) \rightarrow V(\mathcal{C}')$ such that for each block $B \in \mathcal{C}$, the image $\phi(B)$ is a block in \mathcal{C}' . When $|\mathcal{C}| = i$, a configuration \mathcal{C} is called an *i-configuration*. A configuration \mathcal{C} is *even* if for every point a appearing in \mathcal{C} the number $|\{B \mid a \in B \in \mathcal{C}\}|$ of blocks containing a is even.

The notion of minimum distance can be described in the language of combinatorial designs. An $S(2, \mu, v)$ is *r-even-free* if for every integer i satisfying $1 \leq i \leq r$ it contains no even i -configurations. Because the minimum distance of a linear code is the size of a smallest linearly dependent set of columns in its parity-check matrix, the minimum distance of a linear code based on an incidence matrix A of a Steiner 2-design is determined by its even-freeness.

Proposition 3.8: The minimum distance of a linear code whose parity-check matrix forms an incidence matrix of a Steiner 2-design is d if and only if the corresponding Steiner 2-design is $(d - 1)$ -even-free but not d -even-free.

By definition every r -even-free $S(2, \mu, v)$, $r \geq 2$, is also $(r - 1)$ -even-free. If μ is odd, a simple double counting argument shows that an r -even-free $S(2, \mu, v)$ with r even is also $(r + 1)$ -even-free. Because a Steiner 2-design is a linear space in the sense of incidence geometry, every $S(2, \mu, v)$ is trivially μ -even-free.

A nontrivial $S(2, \mu, v)$ may or may not be $(\mu + 1)$ -even-free. For each $\mu \geq 2$, an even $(\mu + 1)$ -configuration that may arise in $S(2, \mu, v)$ s is unique up to isomorphism; they are the dual of the complete graph on $\mu + 1$ vertices. For instance, for the case when $\mu = 3$, up to isomorphism, there exists only one possible even 4-configuration, called the *Pasch* configuration. It can be written by six points and four blocks:

$$\{\{a, b, c\}, \{a, d, e\}, \{f, b, d\}, \{f, c, e\}\}.$$

The unique possible even $(\mu + 1)$ -configurations for $\mu \geq 4$ are sometimes called the *generalized Pasch* configurations in the coding theory literature (see, for example, [46], [52]). Since they are the smallest and unique, an $S(2, \mu, v)$ is $(\mu + 1)$ -even-free if and only if it contains no Pasch configurations for $\mu = 3$ and no generalized Pasch configurations for $\mu \geq 4$.

A fairly tight bound on the maximum even-freeness of an $S(2, 3, v)$ is available.

Theorem 3.9 ([53]): There exists no nontrivial 8-even-free $S(2, 3, v)$.

Hence, by Proposition 3.8, Theorem 3.9, and the fact that every $S(2, 3, v)$ is 3-even-free, we obtain bounds on the minimum distance.

Theorem 3.10: The minimum distance d of a linear code whose parity-check matrix forms an incidence matrix of a nontrivial $S(2, 3, v)$ satisfies the inequalities $4 \leq d \leq 8$.

The problem of avoiding Pasch configurations has long been investigated in various contexts in discrete mathematics. The fundamental question that asks which order v admits an $S(2, 3, v)$ avoiding Pasch configurations was settled in 2000 [54]. Note that such $S(2, 3, v)$ s are 4-even-free and hence are automatically 5-even-free due to their block size being odd number 3.

Theorem 3.11 ([54]): There exists a 5-even-free $S(2, 3, v)$ if and only if $v \equiv 1, 3 \pmod{6}$ except $v = 7, 13$.

While attaining $(\mu + 1)$ -even-freeness in the right portion A of our parity-check matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ is good enough to achieve the goal of reducing the number of codewords of the smallest weight, if one wishes even higher even-freeness, it is required to construct an $S(2, 3, v)$ that simultaneously avoids Pasch and two more even configurations, namely the *grids*

$$\{\{a, b, c\}, \{d, e, f\}, \{g, h, i\}, \{a, d, g\}, \{b, e, h\}, \{c, f, i\}\}$$

and *double triangles*

$$\{\{a, b, c\}, \{a, d, e\}, \{b, f, g\}, \{c, h, e\}, \{d, g, i\}, \{f, h, i\}\}.$$

Unfortunately, while there exist infinitely many $S(2, 3, v)$ s avoiding both Pasch and double triangle configurations [55], no nontrivial examples avoiding grids, let alone 6-even-free $S(2, 3, v)$ s, are known at the time of writing [56]. If a nontrivial $S(2, 3, v)$ that simultaneously avoids the three even configurations exists, it is automatically 7-even-free and hence attains the upper bound given in Theorem 3.10 on the minimum distance of the corresponding linear code.

It is tempting to prove similar theorems on the even-freeness of $S(2, \mu, v)$ s for all $\mu \geq 4$. Unfortunately, while it appears that in principle some of the analogous mathematical arguments likely work [57], it seems very difficult to obtain equally tight bounds and/or complete existence results for relatively high even-freeness for general block size μ . In fact, no nontrivial upper bounds seem to be known on the even-freeness of $S(2, \mu, v)$ s with large μ or, equivalently, on the minimum distances of the corresponding LDPC codes in general. To the best of the authors' knowledge, the only useful and fairly general bound is the one for $S(2, \mu, v)$ s with special automorphisms.

Theorem 3.12 ([58]): If an abelian group acts transitively on the points of a nontrivial r -even-free $S(2, \mu, v)$ with $v > \mu(\mu - 1) + 1$, then $r \leq 2\mu - 1$.

The usefulness of the above bound lies in the fact that the kind of $S(2, \mu, v)$ that is easy to analyze and likely has higher even-freeness than what the trivial lower bound suggests tends to possess the algebraic property considered in Theorem 3.12. In fact, all known nontrivial $S(2, \mu, v)$ s with the highest even-freeness for $\mu \geq 4$ admit such abelian group actions and achieve the upper bound given in Theorem 3.12.

The *affine geometry* $\text{AG}(m, q)$ of dimension m over \mathbb{F}_q is defined as a finite geometry in which the *points* are the vectors in \mathbb{F}_q^m and the *i -dimensional affine subspaces* are the i -dimensional vector subspaces of \mathbb{F}_q^m and their cosets. The points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ form

the points and blocks of an $S(2, q, q^m)$ [44]. Affine geometries provide an explicit construction for nontrivial $S(2, \mu, v)$ s with the highest known even-freeness.

Theorem 3.13 ([59]): For any odd prime power q and positive integer $m \geq 2$ the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ form a $(2q - 1)$ -even-free $S(2, q, q^m)$ which is not $2q$ -even-free.

Affine geometries are not the only known nontrivial $S(2, \mu, v)$ s that attain the upper bound given in Theorem 3.12. The *projective geometry* $\text{PG}(m, q)$ of dimension m over \mathbb{F}_q is a finite geometry whose *points* and *i -dimensional subspaces* are the 1-dimensional vector subspaces and the $(i+1)$ -dimensional vector subspaces of \mathbb{F}_q^{m+1} respectively. The points and 1-dimensional subspaces of $\text{PG}(m, q)$ form the points and blocks of an $S(2, q + 1, \frac{q^{m+1}-1}{q-1})$.

Theorem 3.14 ([28]): For any odd prime power q and positive integer $m \geq 3$ the points and 1-dimensional subspaces of $\text{PG}(m, q)$ form a $(2q + 1)$ -even-free $S(2, q + 1, \frac{q^{m+1}-1}{q-1})$ which is not $(2q + 2)$ -even-free.

Note that because the rank of an incidence matrix of the $S(2, \mu, v)$ from $\text{PG}(m, q)$ with q odd is $v - 1$ [60], the $S(2, q + 1, q^2 + q + 1)$ forming $\text{PG}(2, q)$ with q odd vacuously achieves the highest possible even-freeness $q^2 + q$.

Recently, the first author gave a combinatorial construction for $(\mu + 1)$ -even-free $S(2, \mu, v)$ s [58]. The construction technique recursively combines a $(\mu + 1)$ -even-free $S(2, \mu, v)$ and another $(\mu + 1)$ -even-free $S(2, \mu, w)$ with a particular algebraic property by using a specially designed combinatorial matrix in order to generate a larger $(\mu + 1)$ -even-free $S(2, \mu, vw)$. For the details of the construction, we refer the reader to the original article [58]. As far as the authors are aware, no constructions for $S(2, \mu, v)$ s with even-freeness higher than or equal to $\mu + 1$ are known except the finite geometric and recursive ones.

From the viewpoint of quantum error correction assisted by less noisy qubits, the highly even-free $S(2, \mu, v)$ s based on projective and affine geometries have an additional appealing property. As described in Section II-A, our auxiliary qubits are assumed to be engineered more reliably than the rest so that only phase errors may occur. Hence, it would not be too unnatural to assume that those phase errors that may still occur on these qubits manifest less frequently than bit errors and phase errors on the other qubits. By Equation (1) of Lemma 2.3, in the language of linear codes, this slightly more optimistic assumption translates into the premise that the probability that an error occurs on a fixed check bit, which corresponds to a column of the $(n - k) \times (n - k)$ identity matrix I in $H = \begin{bmatrix} I & A \end{bmatrix}$, is smaller than that on a fixed information bit corresponding to a column of A . The following theorem shows that highly even-free $S(2, \mu, v)$ s from finite geometries can take advantage of this nonuniformity.

Theorem 3.15 ([59], [61]): Let q be an odd prime power and $m \geq 2$ an integer greater than or equal to 2. Define (V, \mathcal{B}) to be the $(2q - 1)$ -even-free $S(2, q, q^m)$ formed by the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$. For any nonempty configuration $\mathcal{C} \subset \mathcal{B}$ of size $|\mathcal{C}| \leq 2q - 1$, it holds that

$$|\mathcal{C}| + \text{odd}(\mathcal{C}) \geq 2q,$$

where $\text{odd}(\mathcal{C})$ is the number of points $v \in V$ contained in an exactly odd number of blocks in \mathcal{C} .

Theorem 3.16 ([61]): Let q be an odd prime power and $m \geq 2$ an integer greater than or equal to 3. Define (V, \mathcal{B}) to be the $(2q+1)$ -even-free $S(2, q+1, \frac{q^{m+1}-1}{q-1})$ formed by the points and 1-dimensional subspaces of $\text{PG}(m, q)$. For any nonempty configuration $\mathcal{C} \subset \mathcal{B}$ of size $|\mathcal{C}| \leq 2q+1$, it holds that

$$|\mathcal{C}| + \text{odd}(\mathcal{C}) \geq 2q+2.$$

The same inequality $|\mathcal{C}| + \text{odd}(\mathcal{C}) \geq 2q+2$ as in Theorem 3.16 holds also for the $S(2, q+1, q^2+q+1)$ from $\text{PG}(2, q)$ with q odd [61]. The point of the above theorems is that a linear code of minimum distance $\mu+1$ defined by a parity-check matrix $H = \begin{bmatrix} I & A \end{bmatrix}$ with A being an incidence matrix of a finite geometric $S(2, \mu, v)$ would perform better if check bits suffer from errors much less likely than information bits. This is because, in a sense, it is effectively of minimum distance 2μ except that there is only one type of small weight codeword, which is the one that consists of one information bit and the corresponding μ check bits. An error of this kind involving check bits would be unlikely to occur if the additional assumption that the more reliable qubits have a sufficiently smaller error probability is valid.

It is notable that this effect of almost doubled minimum distances would be favorable across many different decoding methods and algorithms. In the case of iterative decoding, a nonempty set of bit vertices in a Tanner graph that are not correct after l iterations for all $l \geq l_c$ for some absolute constant l_c is called a *trapping set* [50]. To improve the error floor and slope of the block error rate curve, it is desirable for a parity-check matrix to avoid small trapping sets [62]. While the set of small trapping sets generally varies from algorithm to algorithm and notoriously difficult to identify, codewords of very small weight are surely among them. We will demonstrate the performance of $S(2, \mu, v)$ s based on finite geometries in the context of quantum error correction assisted by more reliable auxiliary qubits in Section IV through simulations.

C. Parity-Check Matrices for Noiseless Qubit Assistance

We now turn our attention to parity-check matrices suitable to entanglement-assisted quantum LDPC codes. In order to put our results in context and show how the two types of assisted quantum error correction are related, we quote the most relevant results that can be found in [28], [31], [33] and also present some useful results that are known in combinatorics but are apparently not found in the quantum coding theory literature. We then give a new method for finding promising high-rate entanglement-assisted quantum LDPC codes at the end of this section.

For entanglement assistance, our parity-check matrices do not need to be in standard form, which was mandatory in the case of quantum error correction assisted by qubits with possible phase errors but no bit errors. The unique requirement in the case of entanglement assistance is that, as mentioned earlier in Section II-B, the 2-rank $\text{rank}(HH^T)$ of the product of our parity-check matrix H and its transpose must be kept small. Thus, in view of Theorem 2.4 and the discussions given

in the previous subsection on error correction performance of incidence matrices of PBDs and their extremely high rates, our interest is in those PBDs that have high even-freeness and very small 2-ranks $\text{rank}(HH^T)$.

To keep our discussion succinct and directly take advantage of most of the material given in the previous subsections, we focus mostly on a class of LDPC codes in which the column weights and the row weights of parity-check matrices are both uniform, that is, *regular* LDPC codes. It is straightforward to see that an incidence matrix of an $S(2, \mu, v)$ is of constant column weight μ and constant row weight $\frac{v-1}{\mu-1}$, providing a parity-check matrix of a regular LDPC code. Since an $S(2, \mu, v)$ contains exactly $\frac{v(v-1)}{\mu(\mu-1)}$ blocks, by Theorem 2.4 and Proposition 3.8, the parameters of the corresponding entanglement-assisted quantum LDPC code are as follows.

Theorem 3.17 ([28]): An incidence matrix H of an r -even-free $S(2, \mu, v)$ that is not $(r+1)$ -even-free gives an entanglement-assisted quantum LDPC code of length $\frac{v(v-1)}{\mu(\mu-1)}$ and dimension $\frac{v(v-1)}{\mu(\mu-1)} - 2\text{rank}(H) + \text{rank}(HH^T)$ that requires $\text{rank}(HH^T)$ ebits for quantum error correction through the LDPC code of the same length, dimension $\frac{v(v-1)}{\mu(\mu-1)} - \text{rank}(H)$, girth 6, and minimum distance $r+1$ formed by H as its parity-check matrix.

As shown in Proposition 3.8, the minimum distance of the LDPC code from an incidence matrix of an $S(2, \mu, v)$ is dictated by its even-freeness. Hence, the bounds and constructions for highly even-free Steiner 2-designs given in Theorems 3.10, 3.11, 3.12, 3.13 and 3.14 are fully and directly applicable here.

Unlike less noisy qubit assistance, however, the dimension of an entanglement-assisted quantum LDPC code depends not only on order v and block size μ but also on 2-ranks concerning the parity-check matrix we use for decoding because, as Theorem 3.17 states, it is $\frac{v(v-1)}{\mu(\mu-1)} - 2\text{rank}(H) + \text{rank}(HH^T)$ for a given incidence matrix H of an $S(2, \mu, v)$. The known results on the possible values of 2-ranks of $S(2, \mu, v)$ s were reviewed in the context of entanglement-assisted quantum LDPC codes in [28], [33]. For convenience, we summarize useful known results here.

The following are the explicit formulas of the 2-ranks of highly even-free projective geometric $S(2, \mu, v)$ s discussed in Section III-B.

Theorem 3.18 ([63]): Let H be an incidence matrix of an $S(2, \mu, v)$ that forms the points and 1-dimensional subspaces of $\text{PG}(m, q)$ with q even. Define $t = \log_2 q$. The 2-rank $\text{rank}(H) = \varphi_e(m, q)$ of the incidence matrix H is given by

$$\varphi_e(m, q) = \sum_{(s_0, s_1, \dots, s_t)} \prod_{j=0}^{t-1} \sum_{i=0}^{L(s_{j+1}, s_j)} l^i \binom{m+1}{i} \binom{m+2s_{j+1}-s_j-2i}{m}$$

where $l = -1$, the sum is taken over all ordered sets (s_0, s_1, \dots, s_t) with $s_0 = s_t$, $s_j \in \mathbb{N}_0$ such that $0 \leq s_j \leq m-1$ and $0 \leq 2s_{j+1}-s_j \leq m+1$ for each $j = 0, \dots, t-1$, and

$$L(s_{j+1}, s_j) = \left\lfloor \frac{2s_{j+1}-s_j}{2} \right\rfloor.$$

Theorem 3.19 ([60]): Let H be an incidence matrix of an $S(2, \mu, v)$ that forms the points and 1-dimensional subspaces of $\text{PG}(m, q)$ with q odd. Then

$$\text{rank}(H) = v - 1 = \frac{q^{m+1} - q}{q - 1}.$$

The 2-rank for the case of a highly even-free Steiner 2-design forming the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ with q even can be expressed by $\varphi_e(m, q)$, that is, the 2-rank of an incidence matrix of an $S(2, \mu, v)$ based on $\text{PG}(m, q)$ with q even.

Theorem 3.20 ([64]): Let H be an incidence matrix of an $S(2, \mu, v)$ that forms the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ with q even. Then the 2-rank of H is given by

$$\text{rank}(H) = \varphi_e(m, q) - \varphi_e(m - 1, q).$$

If q is odd, the 2-rank for the case of $\text{AG}(m, q)$ is full.

Theorem 3.21 ([65]): Let H be an incidence matrix of an $S(2, \mu, v)$ formed by the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ with q odd. Then

$$\text{rank}(H) = v = q^m.$$

If one wishes to employ an $S(2, \mu, v)$ that is not the points and 1-dimensional subspaces of $\text{PG}(m, q)$ or the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$, it is necessary to know the 2-rank of its incidence matrix to compute the dimension through Theorem 3.17. The following are two results on the 2-ranks of $S(2, \mu, v)$ s applicable to half of all general cases.

Theorem 3.22 ([64]): If $\frac{\mu(v-\mu)}{\mu-1}$ is odd, then any incidence matrix H of an $S(2, \mu, v)$ is of full rank, that is, $\text{rank}(H) = v$.

Theorem 3.23 ([64]): If μ is even and $\frac{v-\mu}{\mu-1}$ is odd, then for any incidence matrix H of an $S(2, \mu, v)$, $\text{rank}(H) = v - 1$.

When $\frac{v-\mu}{\mu-1}$ is even, the 2-ranks of incidence matrices of $S(2, \mu, v)$ s may take various values even if v and μ are fixed. In fact, they may vary if Steiner 2-designs are not mutually isomorphic. Hence, if $\frac{v-\mu}{\mu-1}$ is even, finer structural information than the order and block size is needed to calculate the dimension. The most general bounds on the 2-rank of an $S(2, \mu, v)$ read as follows.

Theorem 3.24 ([66]): The 2-rank $\text{rank}(H)$ of an incidence matrix H of an $S(2, \mu, v)$ satisfies inequalities

$$\left\lfloor \frac{1}{2} + \sqrt{\frac{1}{4} + \frac{(v-1)(v-\mu)}{\mu}} \right\rfloor \leq \text{rank}(H) \leq v.$$

The following is very strong theorems for the case when the block size μ is 3.

Theorem 3.25 ([67]): For any $v \equiv 3, 7 \pmod{12}$, where $v = 2^t u - 1$ and u is odd, and any integer i with $1 \leq i < t$, there exists an $S(2, 3, v)$ whose incidence matrix H satisfies the condition that $\text{rank}(H) = v - t + i$.

It is notable that the theorem above covers all orders v to which Theorem 3.22 is not applicable. It is also worth noting that the machinery behind Theorem 3.25 can construct any $S(2, 3, v)$ whose incidence matrix H satisfies the condition that $\text{rank}(H) \leq v - 1$. While the theorem does not treat the

case $\text{rank}(H) = v$, the vast majority of $S(2, 3, v)$ s are actually of full rank. The following theorem provides a simple way to find such Steiner 2-designs.

Theorem 3.26 ([68]): Let H be an incidence matrix of an $S(2, 3, v)$ with a transitive automorphism group. Then $\text{rank}(H) = v$ except when the $S(2, 3, v)$ is the points and 1-dimensional subspaces of $\text{PG}(m, 2)$.

For instance, it is known that for all $v \equiv 1, 3 \pmod{6}$ except for 9 there exists an $S(2, 3, v)$ in which the cyclic group of order v acts regularly on the points [69]. Such an $S(2, 3, v)$ is called *cyclic*. By Theorems 3.26, a cyclic $S(2, 3, v)$ always gives an incidence matrix of full rank except when it is the points and 1-dimensional subspaces of a projective geometry over the binary field \mathbb{F}_2 . The 2-rank of an incidence matrix of an $S(2, 3, v)$ from $\text{PG}(m, 2)$ is known as well.

Theorem 3.27 ([68]): Let H be an incidence matrix of an $S(2, 3, 2^{m+1} - 1)$. Then $\text{rank}(H) \geq 2^{m+1} - m - 2$ with equality if and only if the $S(2, 3, 2^{m+1} - 1)$ is the points and 1-dimensional subspaces of $\text{PG}(m, 2)$.

To compute the dimensions of our entanglement-assisted quantum LDPC codes constructed through Theorem 3.17 with Steiner 2-designs, we also need to know the 2-rank $\text{rank}(HH^T)$ for a given incidence matrix H of an $S(2, \mu, v)$. An important fact to note is that this number $\text{rank}(HH^T)$ is also exactly the number of perfectly noiseless ebits we need to provide. The case when $\text{rank}(HH^T) = 0$ reduces to the case of standard stabilizer codes, where 4-cycles inevitably appear in the Tanner graph of H . Since it is our aim to minimize the number of required ebits, our focus is on those highly even-free $S(2, \mu, v)$ s and similar promising combinatorial designs whose incidence matrices H satisfy the condition that $\text{rank}(HH^T) = 1$.

Recall that an $S(2, \mu, v)$ is a special PBD($v, K, 1$) with $K = \{\mu\}$. The *replication number* r_x of a point $x \in V$ of a PBD (V, \mathcal{B}) is the number of occurrences of x in the blocks of \mathcal{B} . A PBD is *odd-replicate* if for every $x \in V$ the replication number r_x is odd. If the replication number is even for every point, it is *even-replicate*. If $r_x = r_y$ for any two points x and y , we say that the PBD is *equireplicate* (or *regular*) and has replication number r_x . Every $S(2, \mu, v)$ is equireplicate and has replication number $\frac{v-1}{\mu-1}$. Note that while incidence matrices of regular PBDs result in parity-check matrices of *right-regular* LDPC codes in the language of coding theory, their column weights are not necessarily uniform, which means that they may not be *left-regular*. To avoid any confusion, we use the term equireplicate instead of regular when referring to combinatorial designs.

The following is our basic tool to identify combinatorial designs that require as few ebits as possible.

Theorem 3.28 ([33]): Let H be a matrix over \mathbb{F}_2 in which every row and column is of weight greater than 1. $\text{rank}(HH^T) = 1$ if and only if H is an incidence matrix of an odd-replicate PBD($v, K, 1$) in which every point appears more than one block and no block is of size 1.

Assuming that we exclude trifling examples such as LDPC codes of minimum distance 1 or 2, the above theorem essentially says that the number of required ebits is minimized if and only if we use an odd replicate PBDs of index 1. If

we limit ourselves to $S(2, \mu, v)$ s, this means that incidence matrices H of those with $\frac{v-1}{\mu-1}$ odd meet the condition that $\text{rank}(HH^T) = 1$.

While a significant portion of $S(2, \mu, v)$ s including many highly even-free ones given in Section III-B are indeed odd-replicate, not all Steiner 2-designs are. If one wishes to employ even-replicate $S(2, \mu, v)$ s as well while not requiring many ebits, a naive and straightforward way would be to join the identity matrix I to an incidence matrix H to form $H' = \begin{bmatrix} I & H \end{bmatrix}$ as we did for quantum error-correcting codes assisted by less noisy qubits. If we reindex the rows and columns of the extended matrix H' by blocks and points, because the blocks of size 1 corresponding to the columns of I contain no pair of points, H' forms an incidence matrix of an odd-replicate PBD($v, K, 1$). It is easy to verify that the number $\text{rank}(H'H'^T)$ of required ebits becomes 1. The problem of this approach is that the minimum distance of the resulting LDPC code is always $\mu + 1$ regardless of the even-freeness of the $S(2, \mu, v)$ defined by H . Fortunately, because parity-check matrices do not need to be in standard form in entanglement-assistance, there is a simple way around this problem so that one may exploit the promising structure of an incidence matrix of an $S(2, \mu, v)$ even if it is even-replicate.

Theorem 3.29: Let H be an incidence matrix of an even-replicate $S(2, \mu, v)$ and $\mu \geq 2$. Take the $v \times v$ identity matrix I , the v -dimensional all-one vector $J_{1,v} = (1, \dots, 1)$, and the $\frac{v(v-1)}{\mu(\mu-1)}$ -dimensional all-zero vector $\mathbf{0}_{1, \frac{v(v-1)}{\mu(\mu-1)}}$. Define a $(v + 1) \times \left(\frac{v(v-1)}{\mu(\mu-1)} + v\right)$ matrix H' as

$$H' = \begin{bmatrix} I & H \\ J_{1,v} & \mathbf{0}_{1, \frac{v(v-1)}{\mu(\mu-1)}} \end{bmatrix} = \begin{bmatrix} I & H \\ 1 \dots 1 & 0 \dots 0 \end{bmatrix}.$$

H' is an incidence matrix of a PBD($v + 1, K, 1$) such that $\text{rank}(H'H'^T) = 1$. In particular, if the original $S(2, \mu, v)$, (V, \mathcal{B}) , is r -even-free and satisfies the property that for any nonempty configuration $\mathcal{C} \subseteq \mathcal{B}$ with $|\mathcal{C}| \leq r$ and $\text{odd}(\mathcal{C})$ even

$$|\mathcal{C}| + \text{odd}(\mathcal{C}) \geq r + 1, \quad (2)$$

where $\text{odd}(\mathcal{C})$ is the number of points $v \in V$ such that v is contained in an exactly odd number of blocks in \mathcal{C} , then the PBD($v + 1, K, 1$) is r -even-free.

Proof: Take an element $\infty \notin V$ and define a finite set $V' = V \cup \{\infty\}$ of size $|V'| = v + 1$. Index the rows of H' by the elements of V' such that the additional element ∞ is associated with the row $(J_{1,v}, \mathbf{0}_{1, \frac{v(v-1)}{\mu(\mu-1)}})$ and such that the other v rows are associated the same way as in the incidence matrix H . For every unordered pair $\{\infty, v\}$ with $v \in V$, there exists a unique column of H' in which the rows corresponding to ∞ and v both contain 1. Because H is an incidence matrix of a Steiner 2-design, for every unordered pair $\{v, w\}$ such that $v, w \in V$ there exists exactly one column in which the rows indexed by v and w simultaneously contain 1. Hence, the extended matrix H' is an incidence matrix of a PBD($v + 1, K, 1$) with V' as its point set, where the block set \mathcal{B}' consists of v blocks of size 2 and

$\frac{v(v-1)}{\mu(\mu-1)}$ blocks of size μ . It suffices to show that the resulting PBD does not contain any even configurations of size r or smaller if the original $S(2, \mu, v)$ is r -even-free and if every nonempty configuration $\mathcal{C} \subseteq \mathcal{B}$ of size r or smaller such that $\text{odd}(\mathcal{C})$ is even satisfies the inequality $|\mathcal{C}| + \text{odd}(\mathcal{C}) \geq r + 1$. Suppose to the contrary that the PBD (V', \mathcal{B}') contains an even configuration \mathcal{D} of size s for some positive $s \leq r$. Define $\mathcal{D}_H = \{B \in \mathcal{D} \mid B \in \mathcal{B}\}$ to be the set of blocks in \mathcal{D} that are also contained in \mathcal{B} . If $\text{odd}(\mathcal{D}_H)$ is odd, the number of blocks in \mathcal{D} that contain ∞ is odd, contradicting the assumption that \mathcal{D} is an even configuration. If $\text{odd}(\mathcal{D}_H)$ is even, by assumption, $|\mathcal{D}_H| + \text{odd}(\mathcal{D}_H) \geq r + 1$. However, because $\text{odd}(\mathcal{D}_H) = |\mathcal{D} \setminus \mathcal{D}_H|$, we have

$$\begin{aligned} |\mathcal{D}_H| + \text{odd}(\mathcal{D}_H) &= |\mathcal{D}_H| + |\mathcal{D} \setminus \mathcal{D}_H| \\ &= |\mathcal{D}| \\ &= r, \end{aligned}$$

a contradiction. The proof is complete. \blacksquare

Now we illustrate how to apply various theorems presented here and demonstrate how to effectively take advantage of the theorem above through an example case. We first construct through results given in this paper a known class of good entanglement-assisted quantum LDPC codes which originally appeared in [28]. Then we show how Theorem 3.29 extends the class.

By Theorem 3.13, the points and 1-dimensional affine subspaces of $\text{AG}(m, q)$ with q odd form the points and blocks of a $(2q - 1)$ -even-free $S(2, q, q^m)$, which achieves the upper bound on the even-freeness given in Theorem 3.12. Because the replication number of an $S(2, q, q^m)$ is

$$\frac{q^m - 1}{q - 1} = \sum_{i=0}^{m-1} q^i,$$

it is odd-replicate if m is odd, ensuring that the number of required ebits is 1 by Theorem 3.28. Hence, considering also its girth and extremely high rate as an $S(2, \mu, v)$, which was discussed in Section III-B, it would not be too optimistic to expect that an incidence matrix of the $S(2, q, q^m)$ would work well as a parity-check matrix for an entanglement-assisted quantum LDPC code. Applying the rank formula given in Theorem 3.21 to Theorem 3.17, the corresponding entanglement-assisted quantum LDPC code is of length $q^{m-1} \frac{q^m - 1}{q - 1}$ and dimension $q^{m-1} \frac{q^m - 1}{q - 1} - 2q^m$, namely a $[[q^{m-1} \frac{q^m - 1}{q - 1}, q^{m-1} \frac{q^m - 1}{q - 1} - 2q^m]]$ quantum error-correcting code that allows for quantum error correction through the LDPC code of length $q^{m-1} \frac{q^m - 1}{q - 1}$, dimension $q^{m-1} \frac{q^m - 1}{q - 1} - q^m$, girth 6, and minimum distance $2q - 1$ defined by an incidence matrix of the $S(2, q, q^m)$.

The above class of codes based on $\text{AG}(m, q)$ is indeed known to perform in simulations quite similarly to finite geometry LDPC codes proposed in [31] (see [70]). As we have just seen, however, if a large number of ebits are prohibited, this straightforward method only admits $\text{AG}(m, q)$ with m odd when q is also odd. Theorem 3.29 provides a means to exploit the even-replicate $S(2, q, q^m)$ based on $\text{AG}(m, q)$ with m even. In fact, Theorem 3.15 assures that for any m the $(2q - 1)$ -even-free $S(2, q, q^m)$ from $\text{AG}(m, q)$ with q odd

satisfies a stronger condition than Inequality (2) in Theorem 3.29. Thus, its parity-check matrix H can be extended to a $(q^m + 1) \times (q^{m-1} \frac{q^m - 1}{q - 1} + q^m)$ matrix H' which forms an incidence matrix of a $(2q - 1)$ -even-free PBD($q^m + 1, K, 1$). Because $\text{rank}(H'H'^T) = 1$, the entanglement-assisted quantum LDPC code based on this new PBD requires only 1 ebit, as opposed to $q^m - 1$ ebits in the case of the straightforward use of AG(m, q) with m even and q odd (see [28] for the formula for $\text{rank}(HH^T)$ of Steiner 2-designs). Note that the 2-rank of H' , which is required to know to compute the dimension of our code based on the PBD, can be easily obtained. Indeed, it is simply of full rank, that is, $\text{rank}(H') = q^m + 1$. This is because the first q^m rows $\begin{bmatrix} I & H \end{bmatrix}$ are linearly independent due to the identity matrix I and also because no linear combination of these rows coincides with the bottom row

$$\begin{bmatrix} J_{1,q^m} & \mathbf{0}_{1, \frac{q^{m-1}(q^m-1)}{q-1}} \end{bmatrix}$$

due to the fact that q is odd. In fact, all of the first q^m rows must be added together to obtain J_{1,q^m} on the left-hand side while adding up all of them results in the all-one vector on the right-hand side instead of the required zero vector. The fundamental parameters of the new entanglement-assisted quantum LDPC code can be summarized as follows.

Corollary 3.30: For any even integer $m \geq 2$ and odd prime power q there exists an entanglement-assisted quantum error-correcting code of length $q^{m-1} \frac{q^m - 1}{q - 1} + q^m$ and dimension $q^{m-1} \frac{q^m - 1}{q - 1} - q^m - 1$ that requires exactly 1 ebit and can be decoded by the LDPC code of length $q^{m-1} \frac{q^m - 1}{q - 1} + q^m$, dimension $q^{m-1} \frac{q^m - 1}{q - 1} - 1$, girth 6, and minimum distance $2q - 1$.

It is notable that quantum error-correcting codes constructed through Theorem 3.29 are generally expected to have higher dimensions than the original codes used as ingredients because of the extra columns. In the next section we will demonstrate through simulations that this type of quantum LDPC code performs well as expected.

IV. SIMULATION RESULTS

In this section we report simulation results on the performance of our quantum LDPC codes. As noted earlier at the beginning of Section III, we try to be conservative and assume that no exploitable information is available to the receiver regarding possible correlations between bit errors and phase errors. Thus, decoding is done separately for bit errors and phase errors through the sum-product algorithm over two independent binary symmetric channels, where one channel introduces the operator X independently on each physical qubit with probability p_x and the other causes the operator Z to act independently on each physical qubit with probability p_z . Error correction succeeds if the decoder correctly identifies all qubits on which the X operator acted through the sum-product algorithm over one binary symmetric channel and also properly locates all qubits suffering from the Z operator the same way over the other binary symmetric channel.

As in [41], we report the block error rate (BLER) b_p of our LDPC codes over the binary symmetric channel with crossover probability p . Thus, for instance, if one would like a

TABLE I
PARAMETERS OF QUANTUM LDPC CODES

Type ^a	Length	Dimension ^b	Reliable Qubit	Ebit ^c	Rate
RQA	1242	1080	162	0	0.8696
EA	1161	997	0	1	0.8587
H-PEG	1161	997	0	(81)	0.8587
H-AG	1080	918	0	(80)	0.85

^a This column shows if it is assisted by qubits with possible phase errors (RQA), assisted by ebits with no errors (EA), a hypothetical CSS code generated by the PEG algorithm (H-PEG), or a hypothetical one generated by AG(4, 3) (H-AG).

^b The EA code is in catalytic mode for fair comparison. For details, see [72].

^c Parentheses indicate the number of ebits required if Theorem 2.4 were applied.

conservative estimate on the performance of the corresponding quantum LDPC codes over the depolarizing channel with equal error probability $\frac{p}{2}$ for each of the three types of quantum error, the estimated BLER over the quantum channel is $1 - (1 - b_p)^2 \approx 2b_p$. Over a more general Pauli channel with a small error probability p_y for the Pauli operator Y , the same calculation gives a reasonable estimate on the performance of our quantum LDPC codes.

We compare our quantum LDPC codes with *hypothetical* ones that would be available through the CSS construction if there were no constraint on the structure of a parity-check matrix. More specifically, we compete with the ideal situation where any parity-check matrix H of an LDPC code, whether it is in standard form or not, can be used to form a quantum LDPC code regardless of the value of $\text{rank}(HH^T)$. Hence, any parity-check matrix of any linear $[n, k, d]$ code gives rise to a hypothetical $[[n, 2k - n]]$ quantum error-correcting code.

For parity-check matrices of hypothetical codes, we chose good incidence matrices of combinatorial designs found in the coding theory literature and those obtained through the progressive edge-growth (PEG) algorithm, which is among the most successful known methods for designing LDPC codes of relatively short length in the classical domain [71]. Comparison against the structured LDPC codes from promising combinatorial designs makes it easy to see how much the matrix extension processes in Theorems 3.7 and 3.29 affect the performance in our context while the PEG algorithm provides good hypothetical codes for general performance comparison purposes.

Fig. 2 shows the block error rates of our quantum LDPC codes obtained from AG(4, 3) through Theorems 3.7 and 3.29, the hypothetical CSS code based on AG(4, 3), and another hypothetical one generated by the PEG algorithm. The parameters of these codes are summarized in Table I. A close-up of the three based on AG(4, 3) at a small crossover probability region is given in Fig. 3. As shown in these figures and the table, our assisted codes exhibit block error rates comparable to those of the hypothetical ones while significantly reducing the difficulty in implementation and slightly improving information rates.

Fig. 4 compares our quantum LDPC code from AG(3, 5) assisted by less noisy qubits with a hypothetical one constructed by the PEG algorithm. These are both $[[1025, 775]]$ quantum

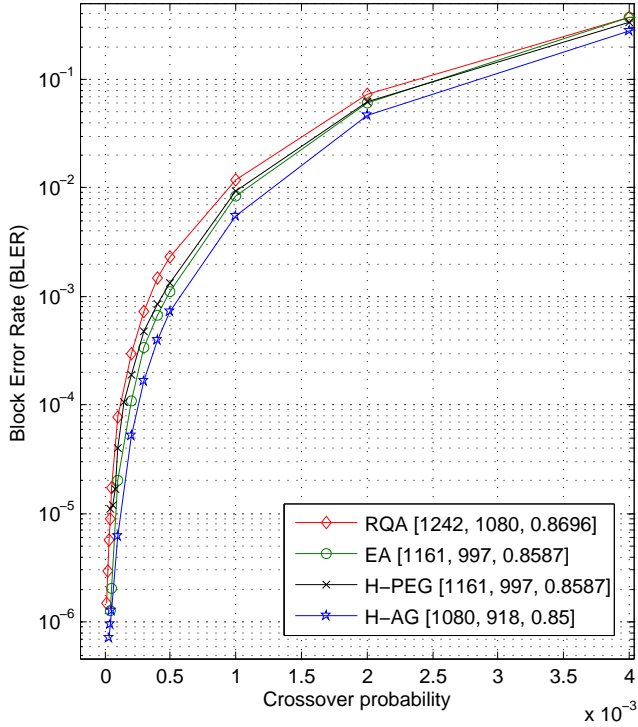


Fig. 2. Block error rates of quantum LDPC codes obtained from $AG(4, 3)$. | RQA and EA refer to assistance by less noisy qubits and ebits respectively. H-PEG and H-AG stand for hypothetical CSS codes based on classical LDPC codes generated by the PEG algorithm and affine geometry $AG(4, 3)$ respectively. The parameters are shown in square brackets in order of [length, dimension, rate].

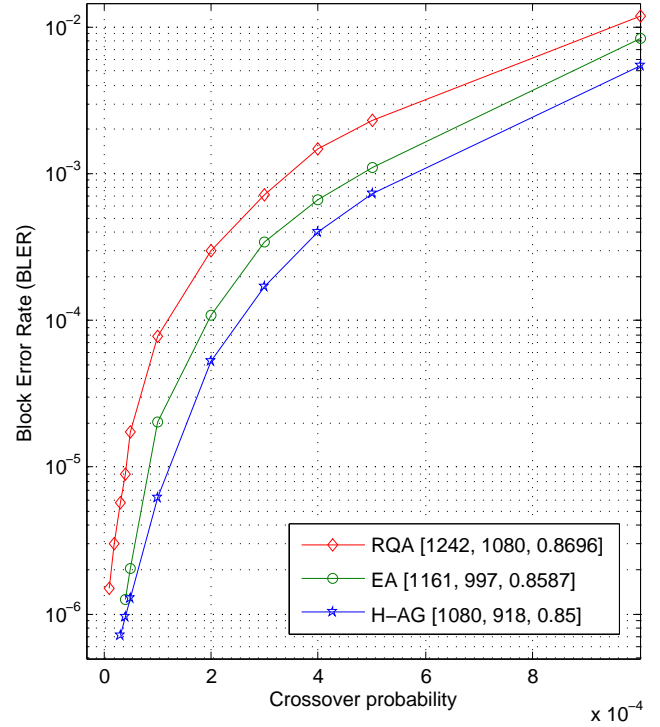


Fig. 3. Block error rates of quantum LDPC codes obtained from $AG(4, 3)$ at lower crossover probabilities. | RQA and EA refer to assistance by less noisy qubits and ebits respectively. H-AG stands for the hypothetical CSS code based on $AG(4, 3)$ that would be available if there were no structural constraint. The parameters are shown in square brackets in order of [length, dimension, rate].

error-correcting codes of rate approximately 0.75. The former requires 250 of all 1025 qubits to be free from bit errors. The latter would need 122 completely error-free qubits on the sender side and another set of 122 perfectly noiseless qubits on the receiver side if maximally entangled pairs were to be used as ebits for quantum error correction. We also plotted block error rates of our code when the 250 less noisy auxiliary qubits experience phase errors less frequently than the rest. This additional assumption can be reasonable because the auxiliary qubits are supposed to be engineered more reliably and protected carefully. As an example, we examined the case when the phase error probability of each auxiliary qubit is a half of that of each noisy one.

Different Steiner 2-designs of the same order and same block size are compared in Fig. 5. Simulation results of LDPC codes from $S(2, 3, 81)$ s that form *Kirkman triple systems* constructed through the *Bose construction* [69] and *5-sparse Steiner triple systems* given in [73] are presented along with the results of those from $AG(4, 3)$. The $S(2, 3, v)$ s from the Bose construction including Kirkman triple systems were studied for use as LDPC codes over additive white Gaussian noise channels in [46], [48]. Incidence matrices of 5-sparse Steiner triple systems are known to avoid small configurations harmful to iterative decoding over a binary erasure channel [55]. While they are not designed specifically for a binary symmetric channel, LDPC codes that are good for these major channels typically perform fairly well, especially if common

harmful structures such as short cycles are avoided. For an analysis of the effects of small configurations in $S(2, 3, v)$ s on iterative decoding, the interested reader is referred to [74]. As expected, our simulation results are, while not identical, overall similar across different $S(2, 3, v)$ s for each type of quantum LDPC code.

V. CONCLUDING REMARKS

We explored the use of quantum error correction assisted by reliable qubits in the context of iterative decoding and demonstrated how one may exploit combinatorics and known designing methods for structured LDPC codes. The range of exploitable classical error-correcting codes for quantum error correction is extended by taking advantage of the fact that some kind of quantum noise is easier to suppress on hardware by physical means. A simple method for creating parity-check matrices of quantum LDPC codes assisted by only one ebit is also given. These codes are shown to have error correction performance comparable to what would be achievable through the same classical ingredients if the CSS construction did not impose any constraints on parity-check matrices.

It should be noted, however, that our results do not imply that the approach presented in this paper removed all difficulties in designing quantum LDPC codes or that the combination of the sum-product algorithm and suitable parity-check matrices is always superior to other coding methods. Rather, Theorem 2.2 and the idea behind Theorem 3.29

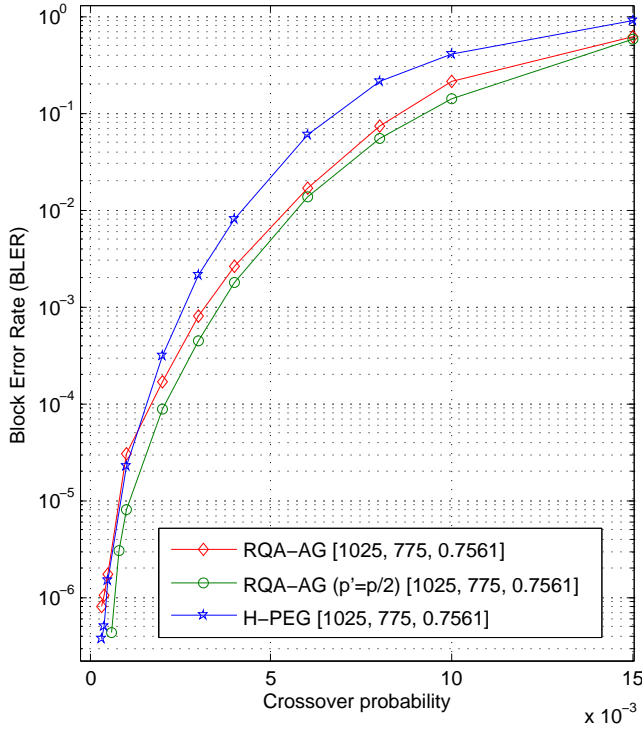


Fig. 4. Comparison between quantum LDPC codes from AG(3, 5) assisted by less noisy qubits and a hypothetical one generated by the PEG algorithm. | RQA-AG refers to the quantum LDPC code based on AG(3, 5) assisted by less noisy qubits. H-PEG refers to a hypothetical one that would be available if the CSS construction did not impose orthogonality on a parity-check matrix. Also plotted are block error rates for when each less noisy qubit of RQA-AG experiences a phase error independently with probability $p' = \frac{p}{2}$, where a phase error occurs on each noisy qubit independently with crossover probability p . The parameters are shown in square brackets in order of [length, dimension, rate].

should be understood as useful tools to circumvent hurdles in designing a variety of quantum error-correcting codes.

To illustrate one limitation in employing LDPC codes for quantum error correction, consider the CSS codes constructed from dual-containing Bose-Chaudhuri-Hocquenghem (BCH) codes (see [75] for the definition of dual-containing BCH codes and their use in quantum error correction). As mentioned in [41], quantum BCH codes outperform all known quantum LDPC codes at rates above 0.8 if the code length is allowed to be several thousand. As far as the authors are aware, the situation does not change if we include all known entanglement-assisted quantum LDPC codes that only require a reasonably small number of ebits. This is partly because it is already not easy to design classical LDPC codes of very high rate that surpass the dual-containing BCH codes in terms of block error rate in the finite length regime even if there is no structural constraint. Hence, if we let less noisy qubits assist quantum error correction, it would still be a challenging task to find a parity-check matrix in standard form that outperforms a dual-containing BCH code of very high rate.

High performance BCH codes are particularly appealing because they also have efficient decoding methods due to their cyclic property (see [76] for decoding in the quantum case). Hence, it would be natural to ask if we can extend the class

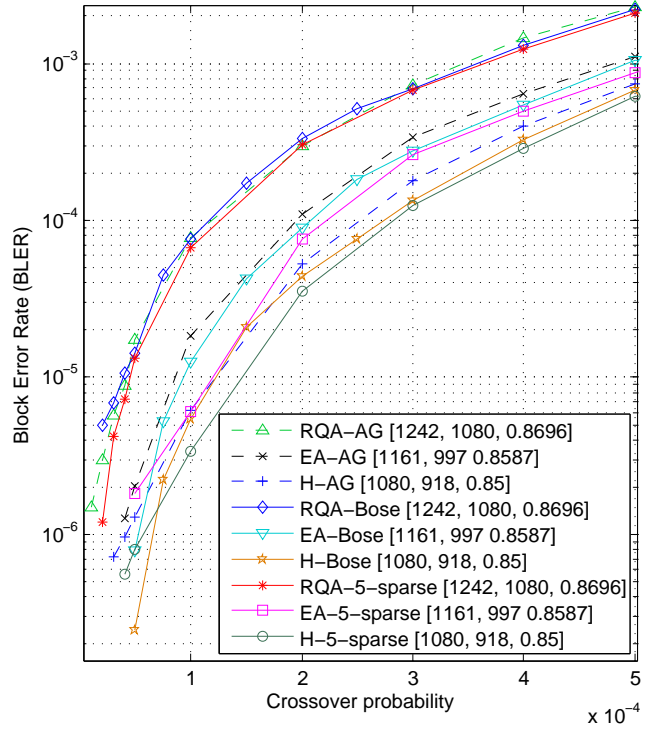


Fig. 5. Block error rates of quantum LDPC codes from different Steiner 2-designs of the same parameters. | RQA and EA refer to assistance by less noisy qubits and ebits respectively. H indicates that a hypothetical situation is assumed where no constraint is imposed on a parity-check matrix by the CSS construction. Codes from AG(4, 3), a Bose-type Kirkman triple system, and a 5-sparse Steiner triple system are shown. All combinatorial designs form $S(2, 3, 81)$ s. The parameters of the corresponding quantum LDPC codes are shown in square brackets in order of [length, dimension, rate].

of quantum BCH codes as we did quantum LDPC codes.

Fortunately, the answer is yes. It is easy to see that Theorem 2.2 we proved in Section II-A is general enough to import BCH codes with their efficient decoding methods and give quantum error-correcting codes with higher rates than the quantum BCH codes through the CSS construction. In fact, as shown in the proof of the theorem, quantum error correction with less noisy qubits extracts the information about quantum noise in the form of a syndrome and exploits it in the same way as in the standard decoding for a linear $[n, k, d]$ code. Hence, in principle, any standard technique that infers errors from syndromes is directly exploitable in the quantum setting as long as the number $2(n - k)$ of less noisy auxiliary qubits falls within the acceptable range. In our case, because the BCH codes in question are of classical information rate $\frac{k}{n} > 0.9$, only a small fraction of qubits need to be free from bit errors. Thus, these BCH codes are ideal classical codes for quantum error correction via less noisy qubits due to their excellent error correction performance and efficient decoding methods.

The slight improvement in quantum information rate with this approach comes from the fact that the standard CSS construction only generates an $[n, 2k - n]$ quantum error-correcting code as it is a special case of Theorem 2.4 when $\text{rank}(HH^T) = 0$. As shown in Theorem 2.2, assistance from less noisy qubits results in a $[[2n - k, k]]$ quantum error-

correcting code. Hence, we always have a slight gain

$$\frac{k}{2n-k} - \frac{2k-n}{n} = \frac{2(n-k)^2}{n(2n-k)} > 0$$

in information rate in the quantum domain.

As we have seen, assistance from reliable qubits seems to have the potential to greatly widen the range of effectively exploitable classical error-correcting codes. We hope that this work stimulates further studies on taking fuller advantage of classical coding theory and also helps find interesting error correction schemes that make use of phenomena unique to the world of quantum information.

REFERENCES

- [1] M. A. Nielsen and I. L. Chuang, *Quantum Computation and Quantum Information*. New York: Cambridge Univ. Press, 2000.
- [2] P. W. Shor, "Scheme for reducing decoherence in quantum computer memory," *Phys. Rev. A*, vol. 52, pp. R2493–R2496, 1995.
- [3] A. M. Steane, "Error-correcting codes in quantum theory," *Phys. Rev. Lett.*, vol. 77, pp. 793–797, 1996.
- [4] D. Gottesman, "Stabilizer codes and quantum error correction," Ph.D. dissertation, California Institute of Technology, 1997.
- [5] A. R. Calderbank, E. M. Rains, P. W. Shor, and N. J. A. Sloane, "Quantum error correction via codes over GF(4)," *IEEE Trans. Inf. Theory*, vol. 44, pp. 1369–1387, 1998.
- [6] A. Cross, G. Smith, J. A. Smolin, and B. Zeng, "Codeword stabilized quantum codes," *IEEE Trans. Inf. Theory*, vol. 55, pp. 433–438, 2009.
- [7] R. Laflamme, C. Miquel, J. P. Paz, and W. H. Zurek, "Perfect quantum error correcting code," *Phys. Rev. Lett.*, vol. 77, pp. 198–201, 1996.
- [8] D. G. Cory et al., "Experimental quantum error correction," *Phys. Rev. Lett.*, vol. 81, pp. 2152–2155, 1998.
- [9] E. Knill, R. Laflamme, R. Martinez, and C. Negrevergne, "Benchmarking quantum computers: the five-qubit error correcting code," *Phys. Rev. Lett.*, vol. 86, pp. 5811–5814, 2001.
- [10] J. Chiaverini et al., "Realization of quantum error correction," *Nature*, vol. 432, pp. 602–605, 2004.
- [11] N. Boulant, L. Viola, E. Fortunato, and D. G. Cory, "Experimental implementation of a concatenated quantum error-correcting code," *Phys. Rev. Lett.*, vol. 94, 130501, 2005.
- [12] C.-Y. Lu et al., "Experimental quantum coding against qubit loss error," *Proc. Natl. Acad. Sci.*, vol. 105, pp. 11 050–11 054, 2008.
- [13] T. Aoki et al., "Quantum error correction beyond qubits," *Nature Physics*, vol. 5, pp. 541–546, 2009.
- [14] P. Schindler et al., "Experimental repetitive quantum error correction," *Science*, vol. 332, pp. 1059–1061, 2011.
- [15] O. Moussa, J. Baugh, C. A. Ryan, and R. Laflamme, "Demonstration of sufficient control for two rounds of quantum error correction in a solid state ensemble quantum information processor," *Phys. Rev. Lett.*, vol. 107, 160501, 2011.
- [16] M. D. Reed et al., "Realization of three-qubit quantum error correction with superconducting circuits," *Nature*, vol. 482, pp. 382–385, 2012.
- [17] X.-C. Yao et al., "Experimental demonstration of topological error correction," *Nature*, vol. 482, pp. 489–494, 2012.
- [18] J. Zhang, R. Laflamme, and D. Suter, "Experimental implementation of encoded logical qubit operations in a perfect quantum error correcting code," *Phys. Rev. Lett.*, vol. 109, 100503, 2012.
- [19] J. Zhang, M. Grassl, B. Zeng, and R. Laflamme, "Experimental implementation of a codeword-stabilized quantum code," *Phys. Rev. A*, vol. 85, 062312, 2012.
- [20] S. Bartz et al., "Demonstrating elements of measurement-based quantum error correction," e-print arXiv:1308.5209, 2013.
- [21] D. Gottesman, "An introduction to quantum error correction and fault-tolerant quantum computation," in *Quantum Information Science and Its Contributions to Mathematics*, ser. Proceedings of Symposia in Applied Mathematics, S. J. Lomonaco, Jr., Ed., vol. 68. Providence, Rhode Island: American Mathematical Society, 2010, pp. 13–58.
- [22] T. A. Brun, I. Devetak, and M.-H. Hsieh, "Correcting quantum errors with entanglement," *Science*, vol. 314, pp. 436–439, 2006.
- [23] M.-H. Hsieh, I. Devetak, and T. A. Brun, "General entanglement-assisted quantum error-correcting codes," *Phys. Rev. A*, vol. 76, 062313, 2007.
- [24] M. M. Wilde and T. A. Brun, "Optimal entanglement formulas for entanglement-assisted quantum coding," *Phys. Rev. A*, vol. 77, 064302, 2008.
- [25] M.-H. Hsieh, T. A. Brun, and I. Devetak, "Entanglement-assisted quantum quasicyclic low-density parity-check codes," *Phys. Rev. A*, vol. 79, 032340, 2009.
- [26] Y. Dong, X. Deng, M. Jiang, Q. Chen, and S. Yu, "Entanglement-enhanced quantum error-correcting codes," *Phys. Rev. A*, vol. 79, 042342, 2009.
- [27] I. B. Djordjevic, "Photonic entanglement-assisted quantum low-density parity-check encoders and decoders," *Optics Lett.*, vol. 35, pp. 1464–1466, 2010.
- [28] Y. Fujiwara, D. Clark, P. Vandendriessche, M. De Boeck, and V. D. Tonchev, "Entanglement-assisted quantum low-density parity-check codes," *Phys. Rev. A*, vol. 82, 042338, 2010.
- [29] M. M. Wilde and T. A. Brun, "Entanglement-assisted quantum convolutional coding," *Phys. Rev. A*, vol. 81, 042333, 2010.
- [30] —, "Quantum convolutional coding with shared entanglement: general structure," *Quantum Inf. Processing*, vol. 9, pp. 509–540, 2010.
- [31] M.-H. Hsieh, W.-T. Yen, and L.-Y. Hsu, "High performance entanglement-assisted quantum LDPC codes need little entanglement," *IEEE Trans. Inf. Theory*, vol. 57, pp. 1761–1769, 2011.
- [32] C.-Y. Lai and T. A. Brun, "Entanglement-assisted quantum error-correcting codes with imperfect ebits," *Phys. Rev. A*, vol. 86, 032319, 2012.
- [33] Y. Fujiwara and V. D. Tonchev, "A characterization of entanglement-assisted quantum low-density parity-check codes," *IEEE Trans. Inf. Theory*, vol. 59, pp. 3347–3353, 2013.
- [34] L. Guo and R. Li, "Linear plotkin bound for entanglement-assisted quantum codes," *Phys. Rev. A*, vol. 87, 032309, 2013.
- [35] Y. Fujiwara, "Quantum error correction via less noisy qubits," *Phys. Rev. Lett.*, vol. 110, 170501, 2013.
- [36] L. Ioffe and M. Mézard, "Asymmetric quantum error-correcting codes," *Phys. Rev. A*, vol. 75, 032345, 2007.
- [37] D. J. C. MacKay, *Information Theory, Inference, and Learning Algorithms*. Cambridge: Cambridge University Press, 2003.
- [38] W. C. Huffman and V. Pless, *Fundamentals of Error-Correcting Codes*. Cambridge: Cambridge Univ. Press, 2003.
- [39] A. R. Calderbank and P. W. Shor, "Good quantum error-correcting codes exist," *Phys. Rev. A*, vol. 54, pp. 1098–1105, 1996.
- [40] M. S. Leifer and D. Poulin, "Quantum graphical models and belief propagation," *Ann. Phys.*, vol. 323, pp. 1899–1946, 2008.
- [41] D. J. C. MacKay, G. Mitchison, and P. L. McFadden, "Sparse-graph codes for quantum error correction," *IEEE Trans. Inf. Theory*, vol. 50, pp. 2315–2330, 2004.
- [42] Y.-J. Wang, B. C. Sanders, B.-M. Bai, and X.-M. Wang, "Enhanced feedback iterative decoding of sparse quantum codes," *IEEE Trans. Inf. Theory*, vol. 58, pp. 1231–1241, 2012.
- [43] M. Denise, J. P. Tillich, and I. Andriyanova, "A family of quantum codes with performances close to the hashing bound under iterative decoding," in *Proc. IEEE Int. Symp. Inf. Theory*, 2013, to appear.
- [44] T. Beth, D. Jungnickel, and H. Lenz, *Design Theory*. Cambridge: Cambridge Univ. Press, 1999.
- [45] R. M. Wilson, "An existence theory for pairwise balanced designs, II: the structure of PBD-closed set and the existence conjecture," *J. Combin. Theory Ser. A*, vol. 13, pp. 246–273, 1972.
- [46] B. Vasić and O. Milenkovic, "Combinatorial constructions of low-density parity-check codes for iterative decoding," *IEEE Trans. Inf. Theory*, vol. 50, pp. 1156–1176, 2004.
- [47] B. Ammar, B. Honary, Y. Kou, J. Xu, and S. Lin, "Construction of low-density parity-check codes based on balanced incomplete block designs," *IEEE Trans. Inf. Theory*, vol. 50, pp. 1257–1268, 2004.
- [48] S. J. Johnson and S. R. Weller, "Resolvable 2-designs for regular low-density parity-check codes," *IEEE Trans. Commun.*, vol. 51, pp. 1413–1419, 2003.
- [49] Y. Kou, S. Lin, and M. P. C. Fossorier, "Low-density parity-check codes based on finite geometries: A rediscovery and new results," *IEEE Trans. Inf. Theory*, vol. 47, pp. 2711–2736, 2001.
- [50] D. V. Nguyen, S. K. Chilappagari, M. W. Marcellin, and B. Vasić, "On the construction of structured LDPC codes free of small trapping sets," *IEEE Trans. Inf. Theory*, vol. 58, pp. 2280–2302, 2012.
- [51] S. J. Johnson, *Iterative Error Correction: Turbo, Low-Density Parity-Check and Repeat-Accumulate Codes*. New York: Cambridge Univ. Press, 2010.
- [52] A. Gruner and M. Huber, "Low-density parity-check codes from transversal designs with improved stopping set distributions," *IEEE Trans. Commun.*, vol. 61, pp. 2190–2200, 2013.

- [53] Y. Fujiwara and C. J. Colbourn, "A combinatorial approach to X-tolerant compaction circuits," *IEEE Trans. Inf. Theory*, vol. 56, pp. 3196–3206, 2010.
- [54] M. J. Grannell, T. S. Griggs, and C. A. Whitehead, "The resolution of the anti-Pasch conjecture," *J. Combin. Des.*, vol. 8, pp. 300–309, 2000.
- [55] C. J. Colbourn and Y. Fujiwara, "Small stopping sets in Steiner triple systems," *Cryptogr. Commun.*, vol. 1, pp. 31–46, 2009.
- [56] Z. Füredi and M. Ruszinkó, "Uniform hypergraphs containing no grids," *Adv. Math.*, vol. 240, pp. 302–324, 2013.
- [57] R. A. Beezer, "Counting configurations in designs," *J. Combin. Theory Ser. A*, vol. 96, pp. 341–357, 2001.
- [58] Y. Fujiwara, "Even-freeness of cyclic 2-designs," e-print arXiv:1210.7516, 2012.
- [59] M. Müller and M. Jimbo, "Erasure-resilient codes from affine spaces," *Discrete Appl. Math.*, vol. 143, pp. 292–297, 2004.
- [60] A. Frumkin and A. Yakir, "Rank of inclusion matrices and modular representation theory," *Israel J. Math.*, vol. 71, no. 3, pp. 309–320, 1990.
- [61] P. Vandendriessche, "On small line sets with few odd points," 2013, preprint.
- [62] M. Ivković, S. K. Chilappagari, and B. Vasić, "Eliminating trapping sets in low-density parity-check codes by using tanner graph covers," *IEEE Trans. Inf. Theory*, vol. 54, pp. 3763–3768, 2008.
- [63] N. Hamada, "The rank of the incidence matrix of points and d -flats in finite geometries," *J. Sci. Hiroshima Univ. Ser. A-I Math*, vol. 32, no. 2, pp. 381–396, 1968.
- [64] —, "On the p -rank of the incidence matrix of a balanced or partially balanced incomplete block design and its applications to error-correcting codes," *Hiroshima Math. J.*, vol. 3, no. 1, pp. 153–226, 1973.
- [65] A. Yakir, "Inclusion matrix of k vs. l affine subspaces and a permutation module of the general affine group," *J. Combin. Theory Ser. A*, vol. 63, pp. 301–317, 1993.
- [66] G. Hillebrandt, "The p -rank of $(0, 1)$ -matrices," *J. Combin. Theory Ser. A*, vol. 60, pp. 131–139, 1992.
- [67] E. F. Assmus, Jr., "On 2-ranks of Steiner triple systems," *Electron. J. Combin.*, vol. 2, R9, 1995.
- [68] J. Doyen, X. Hubaut, and M. Vandensavel, "Ranks of incidence matrices of Steiner triple systems," *Math. Z.*, vol. 163, pp. 251–259, 1978.
- [69] C. J. Colbourn and A. Rosa, *Triple Systems*. Oxford: Oxford Univ. Press, 1999.
- [70] D. C. Clark, "Applications of finite geometries to designs and codes," Ph.D. dissertation, Michigan Technological University, 2011.
- [71] X.-Y. Hu, E. Eleftheriou, and D. M. Arnold, "Regular and irregular progressive edge-growth Tanner graphs," *IEEE Trans. Inf. Theory*, vol. 51, pp. 386–398, 2005.
- [72] T. A. Brun, I. Devetak, and M.-H. Hsieh, "Catalytic quantum error correction," e-print arXiv:quant-ph/0608027, 2006.
- [73] C. J. Colbourn, E. Mendelsohn, A. Rosa, and J. Širáň, "Anti-Mitre Steiner triple systems," *Graphs Combin.*, vol. 10, pp. 215–224, 1994.
- [74] S. Laendner, O. Milenkovic, and J. B. Huber, "Characterization of small trapping sets in LDPC codes from Steiner triple systems," in *Proc. 6th Int. Symp. Turbo Codes Iterative Inf. Processing*, Brest, France, September 2010, pp. 93–97.
- [75] S. A. Aly, A. Klappenecker, and P. K. Sarvepalli, "On quantum and classical BCH codes," *IEEE Trans. Inf. Theory*, vol. 53, pp. 1183–1188, 2007.
- [76] M. Grassl and T. Beth, "Cyclic quantum error-correcting codes and quantum shift registers," *Proc. R. Soc. London Ser. A*, vol. 456, pp. 2689–2706, 2000.